# On the Use of Field RR Lyrae as Galactic Probes. V. Optical and Radial Velocity Curve Templates 

V. F. Braga ${ }^{1,2}$ (1) J. Crestani ${ }^{1,3,4}$ (10) M. Fabrizio ${ }^{1,2}$ (1) G. Bono ${ }^{1,4}$ © C. Sneden ${ }^{5}$ © , G. W. Preston ${ }^{6}$, J. Storm ${ }^{7}$ (D) S. Kamann ${ }^{8}$ (1) , M. Latour ${ }^{9}$ (1) H. Lala $^{10}$, B. Lemasle ${ }^{10}$, Z. Prudil ${ }^{11}$ (1) , G. Altavilla ${ }^{1,2}$ (1), B. Chaboyer ${ }^{12}$ (1) , M. Dall'Ora ${ }^{13}$ (D) , I. Ferraro ${ }^{1}$ (1),   P. B. Stetson ${ }^{20}$ (1), E. Valenti ${ }^{21}$ (1), and M. Zoccali ${ }^{22,23}$ (1)<br>${ }^{1}$ INAF-Osservatorio Astronomico di Roma, via Frascati 33, I-00040 Monte Porzio Catone, Italy<br>${ }^{2}$ Space Science Data Center, via del Politecnico snc, I-00133 Roma, Italy<br>${ }^{3}$ Departamento de Astronomia, Universidade Federal do Rio Grande do Sul, Av. Bento Gonçalves 6500, Porto Alegre 91501-970, Brazil<br>${ }^{4}$ Dipartimento di Fisica, Università di Roma Tor Vergata, via della Ricerca Scientifica 1, I-00133 Roma, Italy<br>${ }^{5}$ Department of Astronomy and McDonald Observatory, The University of Texas, Austin, TX 78712, USA<br>${ }^{6}$ The Observatories of the Carnegie Institution for Science, 813 Santa Barbara St., Pasadena, CA 91101, USA<br>${ }^{7}$ Leibniz-Institut für Astrophysik Potsdam, An der Sternwarte 16, D-14482, Potsdam, Germany<br>${ }^{8}$ Astrophysics Research Institute, Liverpool John Moores University, IC2, Liverpool Science Park, 146 Brownlow Hill, Liverpool,L3 5RF, UK<br>${ }^{9}$ Institute for Astrophysics, Georg-August-University Göttingen, Friedrich-Hund-Platz 1, D-37077 Göttingen, Germany<br>${ }^{10}$ Astronomisches Rechen-Institut, Zentrum für Astronomie der Universität Heidelberg, Mönchhofstr. 12-14, D-69120 Heidelberg, Germany<br>${ }^{11}$ Université de Nice Sophia-antipolis, CNRS, Observatoire de la Côte d'Azur, Laboratoire Lagrange, BP 4229, F-06304 Nice, France<br>${ }^{12}$ Department of Physics and Astronomy, Dartmouth College, Hanover, NH 03755, USA<br>${ }_{14}$ INAF-Osservatorio Astronomico di Capodimonte, Salita Moiariello 16, I-80131 Napoli, Italy<br>${ }^{14}$ Università degli Studi di Napoli Parthenope, Via Amm.F. Acton, 38, I-80133 Napoli, Italy<br>${ }^{15}$ Department of Physics and Astronomy, Iowa State University, Ames, IA 50011, USA<br>${ }^{16}$ Cerro Tololo Inter-American Observatory, NSF's National Optical-Infrared Astronomy Research Laboratory, Casilla 603, La Serena, Chile<br>${ }^{17}$ Instituto de Astrofísica de Canarias, Calle Via Lactea s/n, E-38205 La Laguna, Tenerife, Spain<br>${ }^{18}$ Department of Astronomy, The University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan<br>${ }^{19}$ Department of Physics, Florida Atlantic University, 777 Glades Rd, Boca Raton, FL 33431 USA<br>${ }^{20}$ Herzberg Astronomy and Astrophysics, National Research Council, 5071 West Saanich Road, Victoria, British Columbia V9E 2E7, Canada<br>${ }^{21}$ European Southern Observatory, Karl-Schwarzschild-Str. 2, D-85748 Garching bei Munchen, Germany<br>${ }^{22}$ Instituto Milenio de Astrofísica, Santiago, Chile<br>${ }^{23}$ Pontificia Universidad Catolica de Chile, Instituto de Astrofisica, Av. Vicuña Mackenna 4860, Santiago, Chile Received 2021 May 18; revised 2021 June 20; accepted 2021 June 21; published 2021 September 28


#### Abstract

We collected the largest spectroscopic catalog of RR Lyrae (RRLs) including $\approx 20,000$ high-, medium-, and lowresolution spectra for $\approx 10,000$ RRLs. We provide the analytical forms of radial velocity curve (RVC) templates. These were built using 36 RRLs ( 31 fundamental-split into three period bins-and five first-overtone pulsators) with well-sampled RVCs based on three groups of metallic lines ( $\mathrm{Fe}, \mathrm{Mg}, \mathrm{Na}$ ) and four Balmer lines $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, $\mathrm{H}_{\delta}$ ). We tackled the long-standing problem of the reference epoch to anchor light-curve and RVC templates. For the $V$-band, we found that the residuals of the templates anchored to the phase of the mean magnitude along the rising branch are $\sim 35 \%$ to $\sim 45 \%$ smaller than those anchored to the phase of maximum light. For the RVC, we used two independent reference epochs for metallic and Balmer lines and we verified that the residuals of the RVC templates anchored to the phase of mean RV are from $30 \%$ (metallic lines) up to $45 \%$ (Balmer lines) smaller than those anchored to the phase of minimum RV. We validated our RVC templates by using both the single-point and the three phase point approaches. We found that barycentric velocities based on our RVC templates are two to three times more accurate than those available in the literature. We applied the current RVC templates to Balmer lines RVs of RRLs in the globular NGC 3201 collected with MUSE at VLT. We found the cluster barycentric RV of $V_{\gamma}=496.89 \pm 8.37$ (error) $\pm 3.43$ (standard deviation) $\mathrm{km} \mathrm{s}^{-1}$, which agrees well with literature estimates.


Unified Astronomy Thesaurus concepts: RR Lyrae variable stars (1410); Atomic spectroscopy (2099); Radial velocity (1332); Globular star clusters (656)
Supporting material: machine-readable tables

## 1. Introduction

Pulsating variables are behind numerous breakthroughs in astrophysics. Classical Cepheids (CCs) were used to estimate the distance to M31 and solve the Great Debate concerning the

[^0]extragalactic nature of the so-called Nebulae (Hubble 1926), and to trace, for the first time, the rotation of the Galactic thin disk (Oort 1927; Joy 1939). The size and the age of the universe were revolutionized thanks to the discovery of the difference between CCs, RRLs, and type II Cepheids (TIICs, Baade 1956). Indeed, while they were previously thought to represent the same type of variable stars, it became clear that they represented very distinct populations, with the RRLs and TIICs being very old ( $t \geqslant 10 \mathrm{Gyr}$ ), and the CCs very young ( $t \leqslant 300 \mathrm{Myr}$ ).

Nowadays, CCs are among the most popular calibrators of the extragalactic distance scale (Riess et al. 2019). RRLs, albeit fainter, are excellent standard candles that can provide robust, independent distance measurements even for stellar populations where the young CCs are absent. RRLs obey the welldefined period-luminosity-metallicity relations for wavelengths longer than the $R$-band (Bono et al. 2003; Catelan 2009). As tracers of purely old stellar populations, they can be used to investigate the early formation and evolution of both the Galactic Halo (Fiorentino et al. 2015; Belokurov et al. 2018; Fabrizio et al. 2019) and Bulge (Pietrukowicz et al. 2015; Braga et al. 2018).

It is noteworthy that we lack general consensus on the Galactic Halo structure, in part because different stellar tracers provide different views concerning its spatial structure and the timescale for its formation. Indeed, Carollo et al. (2007) by using Main Sequence (MS), subgiants, and RGs, and Kinman et al. (2012) by using RRLs suggested that the outer halo is more spherical and its density profile is shallower when compared with the inner halo. In contrast, Keller et al. (2008) by using RRLs and Sesar et al. $(2010,2011)$ by using RRLs plus MS stars suggested that the outer halo has a steeper density profile when compared with the inner halo. Deason et al. (2011) by using Blue Horizontal Branch stars and Blue Stragglers found no change in the flattening as a function of the Galactocentric distance (Sesar et al. 2011). More recently, Xue et al. (2015) by adopting a global ellipsoidal stellar density model with Einasto profile found that the models with constant flattening provide a good fit to the entire Halo.

The tension between different measurements may be due to the sample selection of each study. On the one hand, the ages of the RRLs cover a narrow range from $\sim 10$ to $\sim 13$ Gyr. There is evidence that a few RRLs-or stars that mimic RRLs, see Smolec et al. (2013) -are the aftermath of binary evolution, but they only represent a few percent of the populations (Bono et al. 1997; Pietrzyński et al. 2012; Kervella et al. 2019). On the other hand, red giants (RGs) and MS stars, typically used to investigate the Halo, have only very weak age constraints (Conroy et al. 2021). Indeed, all stellar structures less massive than $2 M_{\odot}$ (older than $\sim 0.5-1.0 \mathrm{Gyr}$ ) experience an RG phase, and MS stars also cover a broad range in stellar masses/ages. This means that if the Halo is the result of an intense disruption and merging activity (Monachesi et al. 2019), RG and MS stars are far from being optimal tracers of the early formation, because they are a mixed bag concerning the age distribution.

Field RRLs are less numerous when compared with RG and MS stars, but their narrow age distribution makes them uniquely suited for Galactic archeology. They probe a significant Halo fraction (Galactocentric distance $\leqslant 150 \mathrm{kpc}$ ) with high accuracy. Their individual distances have uncertainties on average smaller than $3 \%-5 \%$, and their accuracy improves when moving from optical to NIR (Longmore et al. 1986; Catelan et al. 2004; Braga et al. 2015). This is a key advantage even in the Gaia era: Gaia EDR3 has an accuracy of $3 \%$ for Halo RRLs ( $G \leqslant 15 \mathrm{mag}$ ) at 1 kpc , and this accuracy will be extended to 2 kpc at the end of the mission (Clementini et al. 2019). RRLs are also valuable targets from the kinematical point of view. In fact, by measuring their velocities, one gets information on the kinematical state of the old population (Halo, Globular Clusters, Bulge). The pioneering work by Layden $(1994,1995)$, based on 302 RRLs, pointed toward a nonsteady formation of the Halo, favoring a
fragmented accretion scenario (Searle \& Zinn 1978). More recently, Zinn et al. (2020) were able to pinpoint the membership of several Halo RRLs to past merger events (Gaia-Enceladus and the Helmi streams; see Helmi et al. 1999, 2018; Myeong et al. 2018). A few Halo RRLs were also associated with the Orphan stream by Prudil et al. (2021), leading to more solid constraints on the origin of the stream itself. Concerning the Bulge, the kinematic properties of RRLs display a duality, with one group of stars associated with the spheroidal component and the other with the Galactic bar (Kunder et al. 2020).

The number of identified RRL is rapidly growing thanks to the enhancements in telescope collecting areas and instrument efficiency. Thanks to long-term optical (Catalina (Drake et al. 2009, 2017), ASAS (Pojmanski 1997), ASAS-SN (Jayasinghe et al. 2019), DES (Stringer et al. 2019), Gaia (Clementini et al. 2019), OGLE (Soszyński et al. 2019), and Pan-STARRS (Sesar et al. 2017)), near-infrared (VVV and VVV-X (Minniti et al. 2011)) and mid-infrared (neo-WISE (Wright et al. 2010)) surveys, more than 200,000 RRLs were identified in the Galactic spheroid. However, RRLs are demanding targets from an observational point of view. Well-sampled time series, meaning at least a dozen properly sampled, photometric measurements, are required for a solid identification and an accurate characterization. The same limitation applies to the measurement of the RRL barycentric radial velocity (RV) $\left(V_{\gamma}\right)$, because it requires multiple measurements to trace the RV variation along the pulsation cycle. To overcome this limitation, several authors have used the radial velocity curve (RVC) of X Ari, observed more than half a century ago by Oke (1966), as a pseudo-template. More recently, RVC templates have been developed for fundamental (RRab) RRLs (Liu 1991; Sesar 2012, henceforth S12). They allow to estimate $V_{\gamma}$ even with a small number of velocity measurements, provided that the $V$-band pulsation properties are known. The current RVC templates are affected by several limitations: despite being based on 22 RRab stars with periods between 0.37 and 0.71 days, the template of Liu (1991) was derived from RVCs with -at most-a few tens of points each. These points are velocities obtained from a heterogeneous set of unidentified metallic lines, since they were collected from several different papers. S12 provided templates for both metallic and $\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}$, and $\mathrm{H}_{\gamma}$ lines with a few hundreds of RV measurements. However, their Balmer templates do not cover the steep decreasing branch, and even more importantly, the templates were based on only six RRab with periods in a very narrow range ( $0.56-0.59$ days). Finally, no RVC templates are available for first-overtone RRLs (RRc).

This work aims at providing new RVC templates for both RRab and RRc variables by addressing all the limitations described above. We adopted a wide set of specific and wellidentified metallic and Balmer lines for both RRab and RRc stars, as well as hundreds of velocity measurement for each template. Because the velocity curves of the RRab display some peculiar variations among themselves, we also separated them into three bins according to their specific shape and pulsation period. Thus, we can provide uniquely precise templates that cover a wide range of intrinsic parameters of these variable stars.

The paper is structured as follows. In Section 2, we investigate the phasing of the optical light curve and discuss on a quantitative basis the difference between the reference
epoch anchored to the luminosity maximum and to the mean magnitude along the rising branch. We present the spectroscopic data set in Section 3, and provide new RVCs and their properties in Section 4. We put together the RVCs and derive the analytical form of the RVC templates in Section 5, discuss the reference epoch to be used to apply the templates in Section 6, and validate them in Section 7. We provide a practical example of how to use the RVC template on spectroscopic observations of NGC 3201 in Section 8. Finally, in Section 9, we summarize the current results and outline future developments of this project.

## 2. Optical Light-curve Templates

Light-curve templates are powerful tools that model the light curve of a periodic variable star. The templates are parameterized with the properties of the variable stars (pulsation mode, period, and amplitude). These come in handy, e.g., to estimate the pulsation properties with a few available data (Stringer et al. 2019), to obtain $O-C$ diagrams that trace the rate of period change (Hajdu et al. 2021), to predict the luminosity of the star at a given phase, and for various other purposes.

The use of both luminosity and RVC templates relies on the use of a reference epoch. This means that the phase zero of the RVC template has to be anchored to a specific feature of the luminosity/RV curve. The most common reference epoch adopted in the field of pulsating variable stars is the time of maximum light in the optical ( $T_{\max }^{\mathrm{opt}}$ ). The RVC templates available in the literature are also anchored to $T_{\max }^{\mathrm{opt}}$ because it matches, within the uncertainties, the time of minimum in the RVC $T_{\text {min }}^{\mathrm{RV}}$. Note that, by "minimum" in the RVC, we mean the numerical minimum, i.e., the epoch of maximum blueshift. This is an approximate choice due to the well-known phase lag between light and RVC (Castor 1971).

Our group introduced a new reference epoch, namely the epoch at which the magnitude along the rising branch of the $V$ band light curve-that is, the section of the light curve where brightness changes from minimum to maximum-becomes equal to the mean $V$-band magnitude ( $T_{\text {mean }}^{\mathrm{opt}}$; Inno et al. 2015; Braga et al. 2019). We thoroughly discussed the advantages of using $T_{\text {mean }}^{\mathrm{opt}}$ versus $T_{\text {max }}^{\mathrm{opt}}$ in the context of NIR light-curve templates for both CCs and RRLs. The reader interested in a detailed discussion is referred to the quoted papers. Here, we summarize the key advantages in adopting $T_{\text {mean }}^{\mathrm{opt}}$ for RRL variables: (i) RRab variables with large amplitudes have RVCs with a "sawtooth" shape, where the maximum can be misidentified by an automatic analytical fit if the phase coverage is not optimal. The rising branch, however, can be more easily fitted. (ii) A significant fraction of RRc variables displays a well-defined bump/dip before the maximum in luminosity. A clear separation between the two maxima is not trivial if the phase coverage is not optimal. (iii) The estimate of $T_{\text {max }}^{\mathrm{opt}}$ is more prone to possible systematics, even with wellsampled light curves, because several RRc and long-period RRab variables display flat-topped light curves i.e., light curves in which the maximum is almost flat for a relatively broad fraction of the phase cycle $(\sim 0.10)$. (iv) $T_{\max }^{\mathrm{opt}}$ is typically estimated either as the top value of the fit of the light cure or the brightest observed point, when the sampling is optimal (e.g., ASAS-SN). This means that $T_{\text {max }}^{\mathrm{opt}}$ is affected by the intrinsic dispersion of the observations and by the time resolution of photometric data. Meanwhile, $T_{\text {mean }}^{\text {opt }}$ is estimated by interpolating the analytical fit the mean magnitude (see Appendix C.1),


Figure 1. Bailey diagram ( $V$-band amplitude vs. logarithmic period) for the RRLs adopted to build the $V$-band light-curve templates. The RRc variables are marked with blue circles and the RRab variables with red squares. The template bins are labeled and the period thresholds are displayed as vertical dashed lines.
which is a very robust property of the star. Therefore, $T_{\text {mean }}^{\mathrm{opt}}$ is intrinsically more robust because its precision is less dependent on sampling.

In the following, we address on a more quantitative basis these key issues in the context of optical light curves. For this purpose, we take advantage of a homogeneous and complete sample of $V$-band light curves for cluster and field RRL variables. In particular, we use visual light curves for RRLs in M4 (Stetson et al. 2014) and in $\omega$ Cen (Braga et al. 2016) together with literature observations for RRLs with BaadeWesselink (BW) analysis (Braga et al. 2019, and references therein). The RRLs in M4 and in $\omega$ Cen have well-sampled light curves, with the number of phase points ranging from hundreds to more than one thousand.

### 2.1. Phasing of Optical Light Curves

We selected 57 RRLs (7 RRc and 50 RRab) from our M4, $\omega$ Cen, and BW RRLs, and separated them into four period bins (see Figure 1). Following Braga et al. (2019), the thresholds are the following: "RRc," "RRabl" (RRab with periods shorter than 0.55 days), "RRab2" (RRab with periods between 0.55 and 0.70 days), and "RRab3" (RRab with periods longer than 0.70 days). See Section 4.4 for a more detailed discussion.

We normalized all the light curves by subtracting the mean magnitude and dividing them by their peak-to-peak amplitude $A(V)$, and estimated $T_{\mathrm{mean}}^{\mathrm{opt}}$ and $T_{\max }^{\mathrm{opt}}$ for the entire sample and the individual values are listed in columns 5 and 6 of Table 1. $T_{\text {mean }}^{\mathrm{opt}}$ was estimated as described in Appendix C.1, and $T_{\max }^{\mathrm{opt}}$ was estimated by converting the phase of maximum light of the model light curve into a Heliocentric Julian Date.

We visually inspected all the reference epochs derived in this work (see Figure 2). To overcome thorny problems in the phasing of light curves, we manually selected the best value of $T_{\max }^{\mathrm{opt}}$ as the HJD of the phase point closest to the maximum for the variables where the fit does not follow closely the data around maximum light. In contrast, no manual selection of $T_{\text {mean }}^{\mathrm{opt}}$ was needed, because its estimation is by its very nature based on a more robust approach. We anchored the phases to both $T_{\text {max }}^{\mathrm{opt}}$ and $T_{\text {mean }}^{\mathrm{opt}}$, and we piled up the light curves into four period bins. We ended up with eight cumulative and normalized light curves: four with $\tau_{0}$ anchored to $T_{\max }^{\mathrm{opt}}$ and four anchored to $T_{\text {mean }}^{\mathrm{opt}}$.


Figure 2. Panels (a)-(h): the left panels-(a), (c), (e), and (g)-display, from top to bottom, cumulative and normalized $V$-band light curves for the four different period bins phased by assuming as reference epoch $\tau_{0}=T_{\text {mean }}^{\text {opt }}$. The individual measurements are marked with light gray dots, while the black solid line shows the analytic fit of the light curve template and the vertical dotted line the reference epoch. The right panels-(b), (d), (f), and (h)-display the histogram of the residuals of individual phase points with respect to the analytical fit. The standard deviations of the distribution of the residuals are also labeled. Panels (i)-(p): same as (a)-(h), but the $V$-band light curves were phased by assuming as reference epoch $\tau_{0}=T_{\max }^{\mathrm{opt}}$.

Finally, we adopted the PEGASUS (PEriodic GAuSsian Uniform and Smooth) function (a series of multiple periodic Gaussians; Inno et al. 2015) to fit the cumulative and normalized light curves. The form of the PEGASUS fit is

$$
\begin{equation*}
P(\phi)=A_{0}+\Sigma_{i} A_{i} \exp \left(-\sin \left(\frac{\pi\left(\phi-\phi_{i}\right)}{\sigma_{i}}\right)^{2}\right) \tag{1}
\end{equation*}
$$

where $A_{0}$ and $A_{i}$ are the zero points and the amplitudes of the Gaussians, while $\phi_{i}$ and $\sigma_{i}$ are the centers and the $\sigma$ of the Gaussians.

Figure 2 displays the cumulative and normalized light curves of the four period bins. The black solid lines plotted in the left and right panels show the analytical fits of the cumulative and normalized $V$-band light curves with PEGASUS functions (see Equation (2)). The coefficients of the PEGASUS fits are listed in Table 2. The standard deviations plotted to the right of the light curves (see also the last column in Table 2) bring forth two interesting results: (i) the standard deviations of the light curves phased by using $T_{\text {mean }}^{\mathrm{opt}}$ are systematically smaller than those phased using $T_{\max }^{\mathrm{opt}}$. The difference for the period bins in which the light curves display a cuspy maximum (RRab1 and

RRab2) is $\sim 37 \%$ smaller, but it becomes $\sim 45 \%$ smaller for the RRc and the RRab3 period bins, because they are characterized by flat-topped light curves. (ii) The cumulative light curves for the RRc and RRab3 period bins phased using $T_{\text {max }}^{\text {opt }}$ show offsets along the rising branch. This mismatch could lead to systematic offsets of $\sim 30 \%$ in $\operatorname{Amp}(V)$ adopted to estimate the mean $\langle V\rangle$ magnitude. Meanwhile, the cumulative light curves phased using $T_{\text {mean }}^{\mathrm{opt}}$ overlap better with each other over the entire pulsation cycle. There is one exception: the RRab3 period bin shows a marginal difference across the phases of maximum in luminosity, but the error in the adopted $\operatorname{Amp}(V)$ is on average a factor of two smaller ( $\sim 15 \%$ ) than those obtained by using $T_{\max }^{\mathrm{opt}}$ as anchor.

The current circumstantial evidence, based on the same photometric data, indicates that the use of a reference epoch anchored to the phase of mean magnitude along the rising branch allows a more accurate phasing with respect to the phase of the maximum in luminosity.

### 2.2. Phase Offset between $T_{\max }^{\mathrm{opt}}$ and $T_{\mathrm{mean}}^{\mathrm{opt}}$

We are aware that large photometric surveys-but also smaller projects focused on variable stars-provide, as

Table 2
Coefficients and Standard Deviations of the PEGASUS Fits to the $V$-Band Light-curve Template

| Template Bin | $N$ | $A_{0}$ | $\begin{aligned} & A_{1} \\ & A_{6} \end{aligned}$ | $\begin{aligned} & \phi_{1} \\ & \phi_{6} \end{aligned}$ | $\begin{aligned} & \sigma_{1} \\ & \sigma_{6} \end{aligned}$ | $\begin{aligned} & A_{2} \\ & A_{7} \end{aligned}$ | $\begin{aligned} & \phi_{2} \\ & \phi_{7} \end{aligned}$ | $\begin{aligned} & \sigma_{2} \\ & \sigma_{7} \end{aligned}$ | $\begin{aligned} & A_{3} \\ & A_{8} \end{aligned}$ | $\begin{aligned} & \phi_{3} \\ & \phi_{8} \end{aligned}$ | $\begin{aligned} & \sigma_{3} \\ & \sigma_{8} \end{aligned}$ | $\begin{aligned} & A_{4} \\ & A_{9} \end{aligned}$ | $\begin{aligned} & \phi_{4} \\ & \phi_{9} \end{aligned}$ | $\sigma_{4}$ | $A_{5}$ | $\phi_{5}$ | $\sigma_{5}$ | $\sigma$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\operatorname{RRc}\left(T_{\text {mean }}^{\text {opt }}\right)$ | 8387 | $-0.5307$ | 1.5916 |  | 0.2750 | $-0.9843$ | 0.5681 | 2.4905 | -3.0543 | 0.0136 | 0.2443 | 1.8137 | 0.6700 | 1.1791 | 1.4625 | 0.0107 | 0.2233 |  |
|  |  |  | 0.4535 | -0.1160 | 0.5254 | ... | ... | ... | ... | ... | ... | ... | ... | ... |  |  |  | 0.049 |
| RRab1 ( $T_{\text {mean }}^{\text {opt }}$ ) | 25956 | $-0.0537$ | 0.0241 | 2.4686 | $-0.3268$ | -0.4234 | 0.1043 | 0.5060 | $-0.6608$ | 0.0204 | 0.2749 | 0.2908 | 0.6227 | 0.6791 | 0.3389 | $-0.0153$ | -0.1522 |  |
|  |  |  | 0.9761 | -0.0829 | 0.4306 | -0.4576 | -0.1291 | 0.2920 | ... | ... | ... | ... | ... | ... |  |  |  | 0.035 |
| RRab2 ( $T_{\text {mean }}^{\text {opt }}$ ) | 26029 | $-0.3753$ | 0.5372 | 3.5621 | 0.9021 | -0.4391 | 0.0470 | 0.1680 | $-0.1191$ | 0.0124 | 0.0657 | 0.1553 | 0.7271 | 0.3803 | 0.0485 | 0.0239 | $-0.0506$ |  |
|  |  |  | 0.6073 | -0.0445 | 0.5040 | -0.2475 | 0.0976 | 0.2170 | $-0.2263$ | 0.1454 | 0.3674 | ... | ... | ... |  |  |  | 0.028 |
| $\operatorname{RRab} 3\left(T_{\text {mean }}^{\text {opt }}\right)$ | 9787 | $-2.3494$ | $-0.6802$ | 4.0938 | 0.3433 | $-0.4498$ | 0.2080 | 0.4693 | -0.2342 | 0.0167 | 0.1190 | 0.6038 | 0.5839 | 0.7883 | $-0.0828$ | 0.0854 | 0.1449 |  |
|  |  |  | 2.8333 | 0.0219 | 1.5893 | ... | ... | ... | ... | ... | ... | ... | ... | ... |  |  |  | 0.036 |
| $\operatorname{RRc}\left(T_{\max }^{\text {opt }}\right)$ | 8387 | $-4.7750$ | -0.6122 | 2.8134 | 0.5686 | 0.0666 | 0.6401 | 0.2090 | 3.5226 | 0.0897 | 0.4441 | -3.9216 | 1.0871 | 0.4602 | 5.4369 | -0.3928 | 3.0508 |  |
|  |  |  | ... | ... | ... | ... | ... | ... | ... | ... | ... | ... | ... | ... |  |  |  | 0.069 |
| RRab1 ( $T_{\text {max }}^{\text {opt }}$ ) | 25956 | $-0.4586$ | $-0.5500$ | 1.1767 | 0.4577 | 0.0221 | 0.8182 | $-0.1123$ | $-1.5880$ | 0.9838 | 0.4772 | 0.9760 | 0.9329 | 0.3396 | $-0.2869$ | 1.3225 | 0.4907 |  |
|  |  |  | -0.4888 | -1.0239 | -0.1842 | 1.3798 | 1.0799 | -1.3487 | ... | ... | $\ldots$ | ... | ... | ... |  |  |  | 0.043 |
| RRab2 ( $T_{\text {max }}^{\text {opt }}$ ) | 26029 | 0.4326 | 0.3797 | 0.5056 | 2.2368 | $-1.3584$ | $-0.0486$ | 1.3392 | $-2.8021$ | 0.0117 | 0.2499 | 0.4585 | 0.7517 | 0.7523 | 2.1228 | 0.0195 | 0.2295 |  |
|  |  |  | 0.7873 | -0.0764 | 0.4384 | ... | ... | ... | ... | ... | ... | ... | ... | ... |  |  |  | 0.035 |
| $\operatorname{RRab} 3\left(T_{\max }^{\mathrm{opt}}\right)$ | 9787 | $-1.5356$ | $-0.2273$ | 1.0032 | 0.3180 | 0.1744 | 0.7173 | 0.5384 | $-0.0470$ | $-0.0120$ | 0.0916 | 1.7028 | 0.5273 | 1.7784 | 0.3626 | 0.7872 | 0.3318 |  |
|  |  |  | 0.1686 | $-0.1563$ | 0.1623 | $\cdots$ | $\cdots$ | $\cdots$ | $\cdots$ | ... | ... | ... | $\cdots$ | ... |  |  |  | 0.055 |

(This table is available in machine-readable form.)


Figure 3. $V$-band $\Delta \Phi$ vs. pulsation period for RRc (blue) and RRab (red) variables. The constant offset for RRc variables and the linear relation for RRab variables are displayed as solid lines. The offset and the linear relation are labeled at the top left corner, along with their standard deviations.
reference epoch, $T_{\max }^{\mathrm{opt}}$. To overcome this difficulty and to provide a homogeneous empirical framework, we investigated the phase offset between $T_{\text {mean }}^{\mathrm{opt}}$ and $T_{\max }^{\mathrm{opt}}$. In particular, we defined the phase difference

$$
\Delta \Phi=\frac{\left(T_{\max }^{\mathrm{opt}}-T_{\operatorname{mean}}^{\mathrm{opt}}\right)}{P} \bmod 1
$$

where mod is the remainder operator. For this purpose, we could adopt a larger sample of visual light curves of 291 RRLs (54 RRc and 237 RRab) from large photometric surveys (Gaia, ASAS, ASAS-SN, and Catalina), from our own photometry of globular clusters ( $\omega$ Cen and M4), and from the literature (BW sample; see caption of Table 11). We found that the phase difference shows, as expected, a trend with the pulsation period (see Figure 3).

In particular, the RRab variables show a quite clear linear trend of phase offset with period $(\Delta \Phi=0.043+0.099 \cdot P)$, with an intrinsic dispersion of 0.024 . The standard deviation for RRc variables is larger, but there is no clear sign of a period dependency. Therefore, we assume a constant phase difference ( $\Delta \Phi=0.223 \pm 0.036$ ) for RRc variables. We also investigated a possible correlation of phase offset with metallicity by adopting the estimates recently provided by Crestani et al. (2021a), but we found none.

## 3. RV Database

To provide new RVC templates, we performed a large spectroscopic campaign aimed at providing RV measurements for both field and cluster RRLs. We reduced and analyzed a large sample of high-, medium-, and low-resolution (HR, MR, and $L R$ ) spectra. This mix of proprietary data and data retrieved from public science archives was supplemented with RVCs of RRLs available in the literature.

### 3.1. Spectroscopic Catalog

We have collected the largest spectroscopic data set-both proprietary and public-for RRLs. Preliminary versions of this spectroscopic catalog have already been used in studies focused on chemical abundances (Fabrizio et al. 2019; Crestani et al. 2021a, 2021b) and RV (Bono et al. 2020a). In this investigation, we have added new spectroscopic data and will discuss in detail
the spectra used for RV measurements. Our data set comprises a total of 23,865 spectra for 10,413 RRLs. Figure 4 shows that the distribution of the RRLs is well-spread over the Galactic Halo. The key properties of the spectra (spectral resolution and signal-to-noise ratio), the spectrographs, and the spectroscopic sample are summarized in Table 3.
The HR sample mainly includes spectra collected with the Las Campanas Observatory du Pont echelle spectrograph (du Pont, 6,208 spectra), plus HR spectra collected from ESO telescopes ( 277 from UVES@VLT, 320 from HARPS@ 3.6 m , and 55 from FEROS@2.2 m MPG). We also have 100 HR spectra from SES@STELLA, 81 from HRS@SALT, 10 from HARPSN@TNG, and 34 from HDS@Subaru. We collected MR spectra from both X-Shooter@VLT (121 spectra) and the LAMOST MR survey ( 1271 spectra). Finally, our spectroscopic data set includes LR spectra from the LAMOST (9099 spectra) and from the SDSS-SEGUE ( 6289 spectra) surveys.

## 4. RV Curves

The main aim of this investigation is to provide RVC templates that can be used to provide $V_{\gamma}$ for RRLs from a few random RV measurements based on a wide variety of spectra. For this purpose, we selected a broad range of strong and weak spectroscopic diagnostics.

### 4.1. RV Spectroscopic Diagnostics

The decision to use multiple spectroscopic diagnostics was made because different lines form at different atmospheric layers. As the RRL are pulsating stars, different lines may trace very different kinematics even when observed at the same phase. Therefore, the resulting velocity curves for different lines may have different shapes and amplitudes. Consequently, combining different hydrogen and/or metallic lines for a single velocity determination would blur the fine detail of the velocity curves and decrease the accuracy of the $V_{\gamma}$ estimate. With this in mind, we performed RV measurements separately with the following diagnostics: four Balmer lines $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\left.\mathrm{H}_{\delta}\right)$, the Na doublet (D1 and D2), the Mg I b triplet ( $\mathrm{Mg} \mathrm{b}_{1}, \mathrm{Mg} \mathrm{b}_{2}$, and $\mathrm{Mg} \mathrm{b}_{3}$ ) and a set of Fe and Sr lines (three lines of the Fe I multiplet 43 and a resonant Sr II line; see Moore 1972).
The laboratory wavelengths of the quoted absorption lines are listed in Table 4. Figure 5 displays the regions of the spectrum of four RRLs where the quoted lines are located. The four RRLs were selected in order to have one RRL for each period bin of the RVC template (see Section 4.4). To measure the RVs for the quoted diagnostics, we performed a Lorentzian fit to the absorption lines by using an automated procedure written in IDL. The wavelength range adopted by the fitting algorithm is fixed according to the spectral resolution of the different spectrographs. Typically, we selected a range in wavelength that is ten full widths at half maximum (FWHM) to the left and ten to the right. The FWHM was estimated as FWHM $=2.355 \times \frac{\lambda_{\mathrm{obs}}}{R}$, where $\lambda$ is the wavelength of the diagnostic and $R$ is the spectral resolution.
The median uncertainties of the single RV estimates for the adopted spectroscopic diagnostics and the standard deviations of the different data sets are listed in Table 5. Note that the different data sets have median uncertainties, on average, smaller than $1.5 \mathrm{~km} \mathrm{~s}^{-1}$.


Figure 4. Aitoff projection in Galactic coordinates $(l, b)$ of the RRL spectroscopic data set. Blue and red symbols display RRc and RRab variables. High-, medium-, and low-resolution spectra are marked with large circles, crosses, and small circles, respectively. RRL for which we only have radial velocity measurements from the literature are plotted as diamonds (Baade-Wesselink sample; see Section 4.2). Black stars indicate the RRLs used to build the RVC templates.

Table 3
Key Properties of the Different Spectroscopic Data Sets Adopted in This Investigation

| Instrument | $N_{\text {spectra }}$ | $N_{\text {RRab }}$ | $N_{\text {RRc }}$ | $R$ | S/N |
| :---: | :---: | :---: | :---: | :---: | :---: |
| High-resolution |  |  |  |  |  |
| du Pont | 6208 | 114 | 76 | 35,000 | 40 |
| FEROS@2.2 m | 55 | 3 | 0 | 48,000 | 13 |
| HARPS-N@TNG | 10 | 0 | 4 | 115,000 | 40 |
| HARPS@3.6 m | 320 | 19 | 6 | 115,000 | 10 |
| HRS@SALT | 81 | 64 | 5 | 40,000 | 50 |
| SES@STELLA | 100 | 0 | 8 | 55,000 | 35 |
| HDS@Subaru | 34 | 23 | 2 | 60,000 | 35 |
| UVES@VLT | 277 | 62 | 8 | 34,540-107,200 | 20 |
| Medium-resolution |  |  |  |  |  |
| X-Shooter@VLT | 121 | 16 | 2 | 4300-18,000 | 45 |
| LAMOST-MR | 1271 | 106 | 66 | 7500 | 22 |
| Low-resolution |  |  |  |  |  |
| LAMOST-LR | 9099 | 4275 | 1935 | 2000 | 22 |
| SEGUE-SDSS | 5110 | 2487 | 1197 | 2000 | 21 |
|  |  | Tot |  |  |  |
|  | 23865 | 7070 | 3343 |  |  |

Notes. Each row gives either the spectrograph or the spectroscopic data set (column 1), the total number of spectra (column 2), the number of RRab and RRc variables (column 3 and 4), the typical spectral resolution (column 5), and the typical S/N@3950 £ (column 6).

### 4.2. RVCs from the Literature

To complement our data set, we collected RVCs of RRLs from the literature (Cacciari et al. 1987; Jones et al. 1987a, 1987b,

Table 4
Wavelengths of the Lines Adopted for Radial Velocity Measurements

| Species | Line ID | $\lambda(\AA)$ |
| :--- | :---: | :---: |
|  | Balmer lines |  |
| $\mathrm{H}_{\alpha}$ | $\mathrm{H}_{\alpha}$ | 6562.80 |
| $\mathrm{H}_{\beta}$ | $\mathrm{H}_{\beta}$ | 4861.36 |
| $\mathrm{H}_{\gamma}$ | $\mathrm{H}_{\gamma}$ | 4340.46 |
| $\mathrm{H}_{\delta}$ | $\mathrm{H}_{\delta}$ | 4101.74 |
|  | Fe group |  |
| $\mathrm{Fe}_{\text {I }}$ | Fe 1 | 4045.81 |
| $\mathrm{Fe}_{\mathrm{I}}$ | Fe 2 | 4063.59 |
| $\mathrm{Fe}_{\mathrm{I}}$ | Fe 3 | 4071.74 |
| Sr II | Sr | 4077.71 |
|  | Mg group |  |
| Mg I | Mg b | 5167.32 |
| Mg I | $\mathrm{Mg} \mathrm{b}_{2}$ | 5172.68 |
| Mg I | Mg b | 5183.60 |
|  | Na group |  |
| Na I | D 1 | 5889.95 |
| Na I | D 2 | 5895.92 |

1988a, 1988b; Liu \& Janes 1989; Clementini et al. 1990; Fernley et al. 1990; Skillen et al. 1993a, 1993b). During the 1980s and 1990s, several bright RRLs were observed both photometrically (optical and NIR) and spectroscopically (velocities from metallic lines) to apply the BW method (Baade 1926; Wesselink 1946) in order to obtain accurate distance determinations. Therefore, we label the set of RVCs from these works as the BW sample. Unfortunately, it was not possible to collect the spectra, therefore


Figure 5. From top to bottom, HR spectra collected with du Pont for RRLs adopted in the four bins of the RVC template. The flux units are arbitrary. The Gaia ID (alternative name in parentheses), iron abundance $([\mathrm{Fe} / \mathrm{H}])$, and the pulsation period are labeled. The four portions of the spectrum display, from left to right, the metallic lines, Mg I b triplet, Na doublet, and the Balmer lines. Each line is marked in blue and labeled. All the spectra presented in this figure were taken at phases near one-third of the rising branch of the RV curve and are only minimally affected by nonlinear phenomena. The first row shows the spectrum of the RRc variable CS Eri. Second row: same as the top, but for the RRab HH Pup. Third row: same as the top, but for the RRab RV Oct. Fourth row: same as the top, but for the RRab AT Ser.
we adopted the RV estimates as provided in the quoted papers. Overall, the BW sample includes 2725 RV measurements for 36 RRab and 3 RRc.
Although this data set is inhomogeneous and based on a mix of weak metallic lines, it is extremely useful to complement our own measurements. Some of the works mentioned above included optical light curves, which we used to validate the robustness of the reference epoch used in the phasing of the RVC template.

### 4.3. Estimate of Barycentric Velocities, RV Amplitudes, and Reference Epochs

To derive the analytic form of the RVC templates, it is necessary to know the pulsation period $(P)$, the reference epoch $\left(V_{\gamma}\right)$, and the RV amplitude $(\operatorname{Amp}(\mathrm{RV}))$ of the RRLs with a well-sampled RV curve. The first two are needed to convert epochs into phases, and the latter two are used for the normalization of the RV curve. The normalization is a crucial
step because the RVC templates have to be provided as normalized curves, with zero mean and unit amplitude.

Our data set includes RV measurements for more than 10,000 RRLs, but only 74 of them have a well-sampled pulsation cycle. A good phase coverage is necessary for the determination of the pulsation properties required for the creation of the RVC template. Reference epochs and $\operatorname{Amp}(\mathrm{RV})$ are particularly sensitive to the quality of the pulsation cycle sampling. We neglected all the RRLs displaying a clear Blazhko effect (a modulation of the pulsation amplitude, both in light and in RV) that would introduce a large intrinsic spread in the RVC templates. Because of this exacting quality control, we derived the analytic form of the RVC template using only a subset of three dozen RRL in the spectroscopic template (31 RRab and 5 RRc ). They cover a broad range in pulsation periods ( $0.27-0.84$ days) and iron abundances $(-2.6 \leqslant[\mathrm{Fe} /$ $\mathrm{H}] \leqslant-0.2$ ). We label these stars with the name "Template Sample" (TS) and their properties are listed in Table 6. The

Table 5
Typical Uncertainties in Radial Velocity Measurements for the Adopted Diagnostics in the Different Spectroscopic Data Sets

| Instrument | eRV(Fe) |  | eRV(Mg) |  | $\mathrm{eRV}(\mathrm{Na})$ |  | $\operatorname{eRV}\left(\mathrm{H}_{\alpha}\right)$ |  | $\mathrm{eRV}\left(\mathrm{H}_{\beta}\right)$ |  | $\mathrm{eRV}\left(\mathrm{H}_{\gamma}\right)$ |  | $\mathrm{eRV}\left(\mathrm{H}_{\delta}\right)$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\operatorname{mdn}_{\left(\mathrm{km} \mathrm{~s}^{-1}\right)} \sigma$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  |
| High resolution |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| du Pont | 0.168 | 2.154 | 0.217 | 1.592 | 0.382 | 0.257 | 0.966 | 0.488 | 0.944 | 0.510 | 1.168 | 1.173 | 1.200 | 1.446 |
| FEROS@2.2 m | 0.145 | 0.023 | 0.166 | 0.028 | 0.161 | 0.033 | 1.027 | 0.162 | 0.819 | 0.130 | 1.008 | 0.110 | 0.991 | 0.063 |
| HARPS-N@TNG | 0.094 | 0.032 | 0.133 | 0.030 | 0.160 | 0.033 | 0.938 | 0.119 | 0.923 | 0.090 | 1.010 | 0.088 | 1.037 | 0.110 |
| HARPS@3.6m | 0.180 | 0.054 | 0.203 | 0.059 | 0.160 | 0.154 | 1.238 | 0.246 | 0.897 | 0.558 | 1.106 | 0.648 | 1.045 | 0.613 |
| HRS@SALT | 0.209 | 2.186 | 0.214 | 0.811 | 0.229 | 0.270 | 1.181 | 0.407 | 1.105 | 0.387 | 1.851 | 0.742 | 2.948 | 2.174 |
| SES@STELLA | 0.329 | 1.094 | 0.178 | 0.411 | 0.000 | 1.491 | 0.808 | 3.193 | 0.677 | 0.278 | 0.785 | 1.323 | 0.801 | 1.379 |
| HDS@Subaru | ... | $\ldots$ | 0.185 | 0.028 | 0.200 | 0.071 | 0.790 | 0.192 | $\ldots$ | $\ldots$ | ... | $\ldots$ | $\ldots$ | ... |
| UVES@VLT | 0.167 | 0.041 | 0.196 | 0.047 | 0.219 | 0.075 | 1.089 | 0.344 | 0.912 | 0.252 | 1.106 | 0.218 | 1.045 | 0.358 |
| Medium resolution |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| X-Shooter@VLT | 0.259 | 0.003 | 0.329 | 0.011 | 0.263 | 0.136 | 1.189 | 0.232 | 1.071 | 0.181 | 1.106 | 0.100 | 1.045 | 0.039 |
| LAMOST-MR | 0.259 | 0.058 | 0.329 | 0.046 | 0.375 | 0.129 | 1.494 | 0.410 | 1.239 | 0.221 | 1.106 | 0.168 | 1.045 | 0.095 |
| Low resolution |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| LAMOST-LR | 2.863 | 1.495 | 3.657 | 1.214 | 4.168 | 1.874 | 4.645 | 1.182 | 3.440 | 0.528 | 3.072 | 0.421 | 1.983 | 0.492 |
| SEGUE-SDSS | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | ... | 4.180 | 0.572 | 3.096 | 0.411 | 2.765 | 0.331 | 2.613 | 0.209 |

Note. Medians (mdn) and standard deviations $(\sigma)$ of the uncertainties on the RV measurements for the $\mathrm{Fe}, \mathrm{Mg}, \mathrm{Na}, \mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}$, and $\mathrm{H}_{\delta}$ lines
individual RV measurements for the TS variables are given in Table 7.

The RRLs in the TS have well-covered RVCs for all the adopted spectroscopic diagnostics. The only exception is a cluster star ( $\mathrm{Cl}^{*}$ NGC 6341 SAW V1) that has good RVCs only for Fe and Mg lines. The number of calibrating RRLs adopted in this work is six times larger than the RRL sample adopted by S12. Moreover, S12 only included RRab variables covering a limited range in pulsation periods ( $0.56-0.59$ days).

To estimate $V_{\gamma}, \operatorname{Amp}(\mathrm{RV})$, the epoch of mean velocity on the decreasing branch and the epoch of minimum velocity ( $T_{\text {mean }}^{\mathrm{RV}}$ and $T_{\text {min }}^{\mathrm{RV}}$, both for Fe and $\mathrm{H}_{\beta} \mathrm{RVCs}$ ), we fitted the RVCs with the PLOESS algorithm, as described in Bono et al. (2020a). Then, we derived $V_{\gamma}$ as the average of the fit and Amp (RV) as the difference between the maximum and the minimum of the fit. The estimates of $V_{\gamma}, \operatorname{Amp}(\mathrm{RV})$ and their uncertainties are provided in Tables 8 and 9 . Note that we provide these estimates for both the Balmer lines and for the averaged RVCs of $\mathrm{Fe}, \mathrm{Na}$, and Mg . By using $V_{\gamma}$ and $\mathrm{Amp}(\mathrm{RV})$, we normalized all the RVCs of the TS RRLs and derived the Normalized RVCs (NRVCs).

### 4.4. Period Bins for the RVC Templates

The shape of both light curves and RVCs of RRLs depends not only on the pulsation mode but also on the pulsation period. To improve the accuracy of the RVC templates available in the literature, we provide independent RVC templates for RRc and RRab variables. Moreover, we divide, for the first time, the typical period range covered by RRab variables into three different period bins. This improvement is strongly required by the substantial variation in pulsation amplitudes (roughly a factor of five) when moving from the blue to the red edge of the fundamental instability strip (Bono \& Cassisi 1999) and in the morphology of both light curves and RVCs (Bono et al. 2020a; Braga et al. 2020).
With this in mind, we adopted the same period bins that were used for the NIR light-curve templates in (Braga et al. 2019) and for the optical light-curve templates in Section 2. The
reasons for the selection of these specific thresholds were already discussed in Braga et al. (2019): they are basically to maximize the number of points per bin without disregarding the change of the curve shape with period, and to separate RRLs with/without Blazhko modulations. The same arguments hold for the current investigation, with the additional advantage that, by adopting the same bins, the whole set of RVC templates and NIR light-curve templates are rooted on homologous subsamples of RRL variables. The Bailey diagram and its velocity amplitude counterpart in Figure 6 show their $\operatorname{Amp}(V), \operatorname{Amp}(\mathrm{RV})$, pulsation periods, and adopted period bins.
We provide 28 RVC templates in total, by considering the combination of seven different spectroscopic diagnostics ( $\mathrm{H}_{\alpha}$, $\mathrm{H}_{\beta}, \mathrm{H}_{\gamma}, \mathrm{H}_{\delta}, \mathrm{Mg}, \mathrm{Na}$, and $\mathrm{Fe}+\mathrm{Sr}$; see Table 4) and four period.

### 4.5. The Reference Epochs of the RVC Templates

The photometric data available in the literature for the TS RRLs were not collected close in time with our spectroscopic data. Therefore, we cannot anchor the RVC templates to the photometric reference epochs (e.g., $T_{\text {mean }}^{\mathrm{opt}}$ ). Small period variations and/or random phase shifts might significantly increase the dispersion of the points in the cumulative RVCs. The phase coverage of the TS RRLs is good enough to provide independent estimates of both the pulsation period and of the reference epoch. Moreover, $T_{\text {mean }}^{\mathrm{opt}}$ matches $T_{\text {mean }}^{\mathrm{R} V(\mathrm{Fe})}$ within 5\% of the pulsation cycle (see Section 6 for more details), therefore we can adopt the latter to compute the RVC templates of the metallic lines ( $\mathrm{Fe}, \mathrm{Mg}$, and Na ). This choice allows anyone to adopt $T_{\text {mean }}^{\mathrm{opt}}$ to phase the RV measurements and then use our templates (see Appendix C for detailed instructions). Note that, to compute the RVC templates of the Balmer lines, we use $T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$ because there is a well-defined difference in phase between $T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$ and $T_{\text {mean }}^{\mathrm{RV} V(\mathrm{Fe})}$. To provide a solid proof of our assumptions, we performed the same test discussed in Section 2.1.

We derived $T_{\text {mean }}^{\mathrm{RV}}$ and $T_{\text {min }}^{\mathrm{RV}}$ from the average of the RVCs for both the Fe group lines and the $\mathrm{H}_{\beta}$ line, taken as representative

Table 6
Calibrating RRLs Used to Derive the RVC Templates


Notes. From left to right, the columns list the Gaia EDR3 ID, the alternative ID, pulsation period, mean visual magnitude, visual amplitude, reference epochs ( $T_{\text {mean }}^{\mathrm{RV}}$ and $T_{\min }^{\mathrm{RV}}$, both for Fe and $\mathrm{H}_{\beta} \mathrm{RVCs}$ ), and the iron abundance and its error.
${ }^{\text {a }}$ The iron abundances are all taken from homogeneous metallicity estimates in Crestani et al. (2021a). The only exception is $\mathrm{Cl}^{*}$ NGC 6341 SAW V1, for which we have adopted the abundance published in Kraft \& Ivans (2003).
of the Balmer lines. As expected, $T_{\text {min }}^{\mathrm{RV}}$ matches to first approximation $T_{\max }^{\mathrm{opt}}$. Indeed, the latter was adopted by Liu (1991) and by Sesar (2012) to anchor their RVC templates. The working hypothesis behind this assumption is that the minimum in the RVC of the metallic lines takes place at the same phases at which the RVCs based on the Balmer lines attain their minimum, i.e., that $T_{\min }^{\mathrm{RV}(\mathrm{Fe})}$ matches $T_{\min }^{\mathrm{RV}(\mathrm{H} \beta)}$. However, we checked this assumption and found that the mean difference in phase between the two epochs is $0.036 \pm 0.051$, with the minimum in the RVCs of Fe group
lines leading the $\mathrm{H}_{\beta}$ minimum. Although the difference is consistent with being zero, its standard deviation is not negligible: a systematic phase drift of one-twentieth around the minimum of the RVC might lead to offsets in the estimate of $V_{\gamma}$ on the order of $10 \mathrm{~km} \mathrm{~s}^{-1}$.

We point out that TS RRLs have multiple estimates of $V_{\gamma}$ and of $\operatorname{Amp}(\mathrm{RV})$-three from the metallic lines plus four from individual Balmer lines-but they only have two reference epochs: one for the Fe group lines, representative of the metallic RVCs, and one for $\mathrm{H}_{\beta}$. The individual estimates of the

Table 7
Individual Radial Velocity Measurements for the Calibrating RRLs

| Name | Species $^{\mathrm{a}}$ | HJD <br> (days) | $\mathrm{RV}^{\mathrm{b}}$ <br> $\left(\mathrm{km} \mathrm{s}^{-1}\right)$ | $\mathrm{eRV}^{\mathrm{c}}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | | Instrument |
| :---: |

Notes. From left to right, the different columns list the name of the variable, the spectroscopic diagnostic, heliocentric Julian date, the radial velocity, the error on the radial velocity, and the spectrograph.
${ }^{\mathrm{a}} \mathrm{Fe}, \mathrm{Mg}$, and Na indicate the average of [ $\left.\mathrm{Fe} 1, \mathrm{Fe} 2, \mathrm{Fe} 3, \mathrm{Sr}\right],[\mathrm{Mg} 2, \mathrm{Mg} 3]$, and [ $\mathrm{Na} 1, \mathrm{Na} 2$ ] lines, respectively (see Section 5.1). In contrast, Balmer lines radial velocity measurements are single.
${ }^{\mathrm{b}}$ Velocity plus heliocentric velocity and diurnal velocity correction.
${ }^{\text {c }}$ Uncertainty on the radial velocity measurements. For the Balmer lines, it is the uncertainty from spectroscopic data reduction. For $\mathrm{Fe}, \mathrm{Mg}$, and Na , it is the standard deviation of the RVs from different lines.
(This table is available in its entirety in machine-readable form.)
reference epochs for the two groups of lines are listed in Table 6.

The basic idea is to have the different RVC templates phased at reference epochs originating from similar physical conditions. Weak metallic lines and strong Balmer lines display a well-defined RV gradient, and their RVCs are also affected by a phase shift because the former form at high optical depths and the latter at low optical depths (Liu \& Janes 1990; Carney et al. 1992; Bono et al. 1994). We verified that the reference epochs for the weak metallic lines ( $\mathrm{Fe}, \mathrm{Mg}$, and Na ) are the same within $\sim 3 \%$ of the pulsation cycle, while those for the Balmer lines, anchored to the $\mathrm{H}_{\beta} \mathrm{RVC}$, are the same within $\sim 5 \%$ of the pulsation cycle.
Figures 7 and 8 display the cumulative and normalized RVCs (CNRVCs) based on the Fe group lines and on the $\mathrm{H}_{\beta}$ line, respectively. Data plotted in these figures display two interesting features worth being discussed in detail: (i) the residuals between observations and analytical fits for the CNRVCs based on the Fe group lines and phased using the reference epoch anchored to $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ are systematically smaller than the residuals of the same CNRVCs anchored to $T_{\min }^{\mathrm{RV}(\mathrm{Fe})}$. The difference ranges from $\sim 30 \%$ for TS RRLs in the period bin RRab1 to $\sim 40 \%$ for TS RRLs in the period bin RRab3. (ii) The impact of the two different reference epochs is even more evident for the Balmer lines (Figure 8). Indeed, the difference in the standard deviation ranges from $\sim 15 \%$ in the period bin RRc to $\sim 45 \%$ for the period bin RRab2. Moreover, the CNRVCs for the RRc and the RRab3 period bin show quite clearly that the reference epoch (vertical dotted line) anchored to $\tau_{0}=T_{\text {mean }}^{\mathrm{RV} V(\mathrm{Fe})}$ takes place at phases that are slightly earlier than the actual minimum (see right panels). This difficulty is associated with the shape of both light curves and RVCs, and it causes larger and asymmetrical residuals when compared with the CNRVCs of the same period bin anchored to the mean
magnitude/systemic velocity (see the histograms plotted in the panels to the right of the CNRVCs).

The current circumstantial evidence indicates that RVC templates based on Fe group RV measurements and anchored to $\tau_{0}=T_{\text {mean }}^{\mathrm{RV}}$ provide $V_{\gamma}$ that are on average $\sim 30 \%$ more accurate than the same RVCs anchored to $\tau_{0}=T_{\text {min }}^{\mathrm{RV}}$. The improvement in using $T_{\text {mean }}^{\mathrm{RV}}$ compared with $T_{\text {min }}^{\mathrm{RV}}$ becomes even more relevant in dealing with the RVCs based on Balmer lines. Indeed, uncertainties are smaller by up to a factor of three (see Section 7). This further supports the use of $T_{\text {mean }}^{\mathrm{RV}}$ as the optimal reference epoch to construct RVC templates.

## 5. RVC Templates

Before deriving the RVC templates, two pending issues need to be addressed: are the RVCs for the different lines in the Fe group, in the Mg Ib triplet, and in the Na doublet, within the errors, the same? Are the mean RVCs of these three groups of lines affected by possible systematics?

### 5.1. Mean RVC Templates for the Three Different Groups of Metallic Lines

Our data set is large enough to derive RVC templates for each single absorption line listed in Table 4. However, our goal is to provide RVC templates that can be adopted as widely as possible. Therefore, we aim to provide one RVC template for each of the Balmer lines and one for each of the metallic groups ( $\mathrm{Fe}, \mathrm{Na}$, and Mg ), making a total of seven different sets of RVC templates. This is feasible only if the RVCs of the lines belonging to the same group have, within the errors, the same intrinsic features, i.e., the same shape, amplitude, and phasing.
In principle, the RVC derived with a specific absorption line is different with respect to the one derived from any another line, because different lines may form in different physical conditions of the moving atmosphere. This is quite obvious for the Balmer series, with $\mathrm{Amp}(\mathrm{RV})$ progressively increasing by $\sim 60 \%-70 \%$ from $\mathrm{H}_{\delta}$ to $\mathrm{H}_{\alpha}$ (S12, Bono et al. 2020a). This is the reason why independent RVC templates have to be provided for each Balmer line. However, for the $\mathrm{Fe}, \mathrm{Mg}$, and Na groups, it is not a priori obvious whether different lines of the same group (e.g., Fe 1 and Fe 2 ) display, within the uncertainties, similar $\operatorname{Amp}(\mathrm{RV})$ and RVC shapes. Therefore, we verified whether the RVCs within the $\mathrm{Fe}, \mathrm{Mg}$, and Na groups agree within uncertainties.

To investigate the difference on a quantitative basis, we inspected the residuals of the RVCs of each line with respect to the average RVC of the group. The left panels in Figure 9 display the residuals of the RV measurements based on Fe 1 , $\mathrm{Fe} 2, \mathrm{Fe} 3$, and Sr lines with respect to the average of the four lines. The middle and right panels are the same, but for the two Mg and the two Na lines. Note that we discarded all the Mg b1 RV measurements because this line is blended with an Fe I line $(5167.50 \AA$ ). This iron line can be as strong as the Mg b1 line itself, or even stronger depending on the Mg abundance and on the effective temperature. Therefore, even using high-resolution spectra, the velocity measurements with the Mg b1 line refer to an absorption feature with a center that changes across the pulsation cycle.

Figure 9 displays, both quantitatively and qualitatively, that for Mg and Na , there is no clear trend within the dispersion of the residuals. The maximum absolute offset is vanishingly small, being always smaller than $0.50 \mathrm{~km} \mathrm{~s}^{-1}$, which is also

# Table 8 

Barycentric Radial Velocities and RV Amplitudes Based on Balmer lines

| Name | $\mathrm{H}_{\alpha}$ |  |  |  | $\mathrm{H}_{\beta}$ |  |  |  | $\mathrm{H}_{\gamma}$ |  |  |  | $\mathrm{H}_{\delta}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | $\begin{aligned} & \operatorname{Amp}(\mathrm{RV}) \\ & \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{aligned}$ | eAmp(RV) | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | $\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{\operatorname{Amp}(\mathrm{RV})}$ | eAmp(RV) | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | $\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{\operatorname{Amp}(\mathrm{RV})}$ | eAmp(RV) | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | $\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{\operatorname{Amp}(\mathrm{RV})}$ | eAmp(RV) |
| YZ Cap | -107.17 | 1.48 | 37.09 | 3.16 | -109.37 | 1.56 | 29.56 | 2.90 | -112.82 | 2.09 | 29.29 | 2.84 | -105.05 | 1.71 | 21.98 | 2.47 |
| DR Cap | -1.66 | 1.18 | 30.42 | 2.94 | -3.25 | 1.21 | 25.31 | 2.54 | -1.83 | 1.82 | 25.25 | 2.94 | -4.42 | 1.89 | 26.86 | 4.80 |
| CS Eri | -145.11 | 1.02 | 48.15 | 3.79 | -146.63 | 1.12 | 35.14 | 2.96 | -146.73 | 1.19 | 33.42 | 2.75 | -142.83 | 1.24 | 28.58 | 2.61 |
| MT Tel | 65.64 | 1.10 | 35.65 | 2.26 | 64.23 | 1.15 | 27.37 | 1.94 | 65.26 | 1.21 | 25.80 | 1.91 | 66.87 | 1.29 | 22.03 | 1.81 |
| SV Scl | -14.60 | 1.11 | 42.70 | 3.19 | -15.94 | 1.21 | 32.64 | 2.73 | -14.93 | 1.34 | 29.18 | 2.58 | -11.09 | 1.56 | 23.17 | 2.48 |
| AV Peg | -57.29 | 0.94 | 92.66 | 4.07 | -62.00 | 0.98 | 72.66 | 3.27 | -66.26 | 1.59 | 67.53 | 3.21 | -59.05 | 1.55 | 59.07 | 2.98 |
| HH Pup | 18.11 | 1.02 | 110.28 | 4.30 | 17.05 | 1.02 | 86.75 | 3.74 | 15.65 | 1.33 | 80.19 | 3.29 | 20.66 | 1.21 | 70.89 | 2.98 |
| V0445 Oph | -19.51 | 1.27 | 90.33 | 7.07 | -24.41 | 1.18 | 68.82 | 5.16 | -28.84 | 1.53 | 63.08 | 4.79 | -21.29 | 1.62 | 58.15 | 4.98 |
| ST Vir | -1.26 | 1.40 | 95.80 | 5.42 | -4.63 | 1.54 | 73.00 | 4.24 | -2.24 | 1.61 | 75.59 | 4.48 | -3.03 | 1.57 | 61.47 | 3.81 |
| W Crt | 60.86 | 1.21 | 101.33 | 4.17 | 59.32 | 1.15 | 84.23 | 3.55 | 56.66 | 1.57 | 76.17 | 3.28 | 61.81 | 1.38 | 66.79 | 2.89 |
| VX Her | -376.05 | 1.72 | 105.14 | 7.29 | -376.92 | 1.85 | 80.88 | 6.88 | -379.06 | 2.35 | 75.14 | 5.33 | -373.62 | 2.01 | 57.18 | 4.06 |
| SW Aqr | -49.20 | 1.08 | 103.54 | 4.13 | -49.17 | 1.13 | 76.46 | 3.82 | -49.25 | 1.36 | 80.24 | 3.42 | -46.23 | 1.55 | 66.89 | 2.99 |
| DX Del | -56.88 | 0.95 | 92.90 | 5.42 | -60.81 | 0.98 | 71.37 | 3.92 | -62.74 | 1.31 | 66.59 | 3.78 | -58.20 | 1.19 | 58.22 | 3.24 |
| UU Vir | -12.08 | 0.89 | 110.49 | 6.31 | -12.51 | 0.93 | 94.20 | 5.79 | -11.12 | 1.45 | 88.31 | 4.28 | -9.62 | 1.28 | 75.87 | 3.53 |
| V0440 Sgr | -64.10 | 1.67 | 98.91 | 6.24 | -64.22 | 1.93 | 84.25 | 6.78 | -60.83 | 1.96 | 74.25 | 5.63 | -59.86 | 1.99 | 67.85 | 5.09 |
| ST Leo | 165.14 | 1.72 | 114.08 | 8.18 | 165.08 | 1.05 | 82.29 | 5.35 | 163.93 | 1.55 | 86.00 | 5.88 | 169.63 | 1.68 | 69.03 | 4.78 |
| $V$ Ind | 200.88 | 0.95 | 101.31 | 3.39 | 200.31 | 0.85 | 78.22 | 2.58 | 201.52 | 1.04 | 73.88 | 2.48 | 203.00 | 0.95 | 61.37 | 2.07 |
| AN Ser | -41.71 | 1.35 | 95.50 | 6.29 | -44.51 | 1.50 | 73.11 | 4.87 | -51.79 | 1.77 | 69.88 | 4.65 | -42.15 | 1.55 | 57.87 | 3.77 |
| RR Cet | -77.69 | 1.06 | 107.53 | 5.23 | -77.19 | 1.28 | 82.88 | 4.13 | -75.66 | 1.72 | 77.60 | 4.02 | -73.85 | 1.42 | 72.01 | 3.67 |
| DT Hya | 75.15 | 1.30 | 110.59 | 5.82 | 77.66 | 1.12 | 82.92 | 4.23 | 77.45 | 1.43 | 81.03 | 4.31 | 82.13 | 1.46 | 71.75 | 3.87 |
| TY Gru | -12.28 | 1.54 | 105.82 | 4.14 | -11.13 | 1.69 | 80.50 | 3.24 | -6.89 | 1.95 | 81.11 | 3.37 | -8.79 | 2.19 | 63.10 | 3.13 |
| RV Oct | 137.23 | 1.42 | 104.50 | 3.63 | 139.60 | 1.15 | 85.93 | 3.00 | 141.51 | 1.35 | 85.52 | 3.12 | 144.78 | 1.29 | 75.42 | 2.86 |
| CD Vel | 239.70 | 0.95 | 96.88 | 3.86 | 240.15 | 0.83 | 77.77 | 3.02 | 241.53 | 1.11 | 74.63 | 2.95 | 243.02 | 1.07 | 62.65 | 2.52 |
| WY Ant | 201.31 | 0.93 | 108.83 | 5.62 | 202.26 | 0.78 | 84.27 | 4.27 | 204.58 | 1.00 | 76.90 | 4.12 | 206.15 | 0.99 | 67.96 | 3.45 |
| BS Aps | -106.76 | 1.25 | 85.20 | 3.01 | -108.88 | 1.08 | 65.30 | 2.28 | -105.39 | 1.22 | 61.26 | 2.21 | -103.51 | 1.36 | 56.18 | 2.22 |
| Z Mic | -58.64 | 0.94 | 92.91 | 3.61 | -61.08 | 0.79 | 74.78 | 2.88 | -58.05 | 0.98 | 68.05 | 2.67 | -57.44 | 1.05 | 60.64 | 2.41 |
| XZ Aps | 192.39 | 1.02 | 104.24 | 3.14 | 195.55 | 0.91 | 85.62 | 2.62 | 197.67 | 1.46 | 85.91 | 2.85 | 198.92 | 1.60 | 67.22 | 2.49 |
| SX For | 243.01 | 1.25 | 91.87 | 5.73 | 242.72 | 0.99 | 69.79 | 3.98 | 244.01 | 1.17 | 65.42 | 3.86 | 245.87 | 1.01 | 59.78 | 3.45 |
| SS Leo | 160.70 | 0.99 | 103.73 | 5.54 | 161.11 | 1.17 | 83.63 | 4.13 | 164.42 | 1.56 | 80.95 | 4.00 | 164.07 | 1.61 | 73.69 | 3.62 |
| DN Aqr | -229.89 | 0.87 | 101.94 | 5.14 | -228.68 | 0.84 | 83.05 | 4.38 | -229.00 | 1.04 | 81.96 | 4.14 | -225.74 | 1.31 | 64.72 | 3.50 |
| W Tuc | 56.91 | 1.07 | 114.34 | 5.36 | 59.96 | 0.99 | 85.68 | 4.10 | 63.23 | 1.20 | 82.03 | 3.88 | 65.48 | 1.15 | 69.92 | 3.41 |
| X Ari | -41.58 | 1.08 | 109.64 | 4.40 | -39.73 | 0.97 | 89.55 | 3.47 | -37.97 | 1.27 | 77.70 | 3.10 | -36.22 | 1.43 | 67.65 | 2.86 |
| VY Ser | -147.44 | 0.95 | 97.31 | 10.27 | -148.75 | 0.81 | 79.42 | 6.68 | -146.66 | 0.94 | 70.92 | 5.80 | -146.87 | 0.87 | 64.06 | 4.70 |
| AT Ser | -71.38 | 1.00 | 102.45 | 7.35 | -70.38 | 1.16 | 79.59 | 5.50 | -68.68 | 1.42 | 79.90 | 5.60 | -63.33 | 1.44 | 68.66 | 4.63 |
| V0384 Tel | 302.85 | 1.45 | 87.45 | 10.13 | 303.03 | 0.98 | 71.76 | 9.27 | 303.04 | 1.99 | 60.79 | 8.15 | 303.90 | 1.75 | 55.80 | 6.81 |

Note. Cl* NGC 6341 SAW V1 does not appear because the Balmer RV measurements for this star are not accurate.
(This table is available in machine-readable form.)

Table 9
Barycentric Radial Velocities and RV Amplitudes Based on Metallic Lines

| Name | Fe |  |  |  | Mg |  |  |  | Na |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | $\begin{gathered} \text { Amp } \\ \text { (RV) } \\ \mathrm{Rs}^{-1} \text { ) } \end{gathered}$ | eAmp (RV) | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | Amp <br> (RV) $\left.\mathrm{s}^{-1}\right)$ | eAmp (RV) | $V_{\gamma}$ | $\mathrm{e} V_{\gamma}$ | Amp (RV) $\left.\mathrm{s}^{-1}\right)$ | eAmp (RV) |
| YZ Cap | -109.60 | 0.16 | 26.34 | 1.09 | -109.05 | 0.16 | 26.47 | 1.55 | -109.05 | 0.41 | 27.76 | 2.33 |
| DR Cap | -3.34 | 0.18 | 22.77 | 1.17 | -2.19 | 0.20 | 23.00 | 1.72 | -3.25 | 0.48 | 27.30 | 2.76 |
| CS Eri | -145.25 | 0.12 | 29.08 | 1.21 | -145.31 | 0.17 | 29.17 | 1.74 | -145.60 | 0.39 | 31.43 | 2.61 |
| MT Tel | 66.43 | 0.14 | 22.82 | 0.73 | 65.98 | 1.17 | 22.99 | 1.90 | 65.32 | 0.71 | 23.19 | 1.53 |
| SV Scl | -14.14 | 0.19 | 26.87 | 1.14 | -13.98 | 0.17 | 26.80 | 1.56 | -15.68 | 0.30 | 26.81 | 2.06 |
| AV Peg | -60.48 | 0.28 | 63.52 | 1.41 | -56.63 | 0.32 | 64.58 | 2.00 | -55.89 | 0.69 | 68.71 | 3.06 |
| HH Pup | 18.59 | 0.21 | 69.27 | 1.36 | 19.65 | 0.29 | 71.16 | 1.95 | 19.00 | 0.65 | 71.42 | 2.82 |
| V0445 Oph | -21.96 | 0.32 | 57.42 | 2.25 | -18.57 | 0.44 | 57.16 | 2.98 | -18.36 | 0.67 | 61.45 | 4.70 |
| ST Vir | -4.52 | 0.27 | 60.60 | 1.74 | -1.43 | 0.54 | 62.44 | 2.62 | -2.50 | 0.64 | 66.75 | 3.83 |
| W Crt | 61.38 | 0.28 | 67.49 | 1.27 | 63.66 | 0.50 | 68.39 | 1.84 | 62.69 | 0.70 | 70.41 | 2.76 |
| VX Her | -375.68 | 0.50 | 61.93 | 1.86 | -374.55 | 0.65 | 62.76 | 2.62 | -375.59 | 0.98 | 62.61 | 3.68 |
| SW Aqr | -48.64 | 0.23 | 62.54 | 1.25 | -47.21 | 0.30 | 62.31 | 1.72 | -47.80 | 0.48 | 64.74 | 1.90 |
| DX Del | -58.96 | 0.24 | 53.90 | 1.48 | -55.69 | 0.54 | 54.27 | 2.05 | -55.59 | 0.60 | 57.46 | 3.39 |
| UU Vir | -11.15 | 0.20 | 67.45 | 1.46 | -3.30 | 0.45 | 69.21 | 2.11 | -8.80 | 0.62 | 69.22 | 3.07 |
| V0440 Sgr | -60.99 | 0.48 | 63.10 | 2.60 | ... | ... | ... | ... | -36.16 | 1.64 | 95.08 | 9.03 |
| ST Leo | 165.87 | 0.29 | 65.39 | 2.12 | 168.07 | 0.34 | 65.51 | 2.94 | 167.40 | 0.54 | 66.96 | 4.44 |
| V Ind | 201.04 | 0.16 | 56.09 | 0.83 | 201.86 | 0.25 | 56.48 | 1.24 | 201.17 | 0.42 | 57.87 | 1.39 |
| AN Ser | -43.34 | 0.33 | 59.90 | 1.84 | -39.68 | 0.54 | 60.80 | 2.56 | -38.94 | 1.01 | 60.08 | 3.96 |
| RR Cet | -75.49 | 0.18 | 61.77 | 1.62 | -74.13 | 0.37 | 61.42 | 2.17 | -74.39 | 0.55 | 63.03 | 3.38 |
| DT Hya | 80.03 | 0.30 | 63.78 | 1.64 | 82.12 | 0.41 | 63.26 | 2.34 | 81.90 | 0.68 | 61.90 | 3.09 |
| TY Gru | -7.89 | 0.41 | 58.61 | 1.05 | -6.70 | 0.43 | 61.07 | 1.50 | -7.26 | 0.55 | 61.90 | 2.15 |
| RV Oct | 141.86 | 0.32 | 65.28 | 1.10 | 142.46 | 0.45 | 65.63 | 1.58 | 142.38 | 1.01 | 69.55 | 2.67 |
| CD Vel | 240.96 | 0.17 | 52.13 | 0.93 | 241.52 | 0.22 | 53.07 | 1.32 | 240.73 | 0.43 | 52.77 | 1.88 |
| WY Ant | 204.31 | 0.17 | 59.80 | 1.42 | 205.28 | 0.25 | 59.45 | 2.01 | 204.40 | 0.45 | 64.67 | 3.26 |
| BS Aps | -106.68 | 0.26 | 50.42 | 0.82 | -105.79 | 0.46 | 49.27 | 1.24 | -106.61 | 0.59 | 49.51 | 1.67 |
| Z Mic | -59.52 | 0.18 | 52.94 | 0.98 | -57.99 | 0.42 | 52.71 | 1.45 | -58.17 | 0.56 | 53.89 | 2.06 |
| XZ Aps | 198.49 | 0.24 | 63.52 | 0.92 | 198.27 | 0.29 | 63.02 | 1.28 | 198.34 | 0.50 | 63.28 | 1.88 |
| SX For | 244.48 | 0.27 | 51.57 | 1.44 | 245.34 | 0.38 | 52.53 | 2.10 | 245.31 | 0.47 | 51.15 | 3.01 |
| SS Leo | 162.46 | 0.23 | 61.51 | 1.36 | 164.25 | 0.35 | 61.74 | 1.99 | 163.50 | 0.37 | 61.84 | 1.96 |
| DN Aqr | -226.83 | 0.20 | 54.58 | 1.39 | -226.24 | 0.46 | 52.91 | 1.96 | -227.23 | 0.50 | 57.35 | 2.51 |
| W Tuc | 64.98 | 0.19 | 63.66 | 1.35 | 64.62 | 0.27 | 63.19 | 1.89 | 64.54 | 0.75 | 65.87 | 3.20 |
| X Ari | -37.37 | 0.22 | 57.66 | 1.06 | -36.70 | 0.51 | 57.26 | 1.61 | -36.94 | 0.76 | 56.32 | 2.24 |
| $\begin{gathered} \mathrm{Cl}^{*} \text { NGC } 6341 \\ \text { SAW V1 } \end{gathered}$ | -125.84 | 1.13 | 56.24 | 4.54 | ... | .. | $\ldots$ | $\ldots$ | ... | ... | ... | $\ldots$ |
| VY Ser | -147.11 | 0.27 | 51.30 | 1.85 | -145.15 | 0.27 | 48.82 | 2.46 | -145.45 | 0.43 | 50.41 | 3.92 |
| AT Ser | -67.66 | 0.23 | 60.01 | 2.57 | -66.03 | 0.36 | 56.55 | 2.98 | -65.78 | 0.57 | 62.87 | 5.45 |
| V0384 Tel | 305.34 | 0.29 | 42.85 | 3.24 | 304.55 | 0.43 | 41.72 | 4.21 | 305.29 | 0.63 | 46.36 | 6.17 |

(This table is available in machine-readable form.)
smaller than our RV uncertainty. This means that these lines trace the dynamics of the same atmospheric layer and they can be averaged in order to derive a single set of RVC templates for both Mg and Na groups.

The same outcome does not apply to Fe and Sr lines. Indeed, although the average offsets are all smaller than the dispersions, they seem to follow a trend. More specifically, the average offset of the Fe 1 curve is always smaller than the other, while the average offset of the Sr curve is typically larger. This is mostly due to the interplay of a small difference in $\operatorname{Amp}(\mathrm{RV})$ (generally increasing from Fe 1 to Sr ) and a mild trend in the average velocity (generally increasing from Fe 1 to Sr ). However, all these offsets are smaller than the dispersions $\left(\leqslant 2.5 \mathrm{~km} \mathrm{~s}^{-1}\right)$ and similar to the intrinsic dispersion of the RVC templates (see Section 5.2). We also note that the standard deviation of the points around the offsets is, within the uncertainties, constant along the pulsation cycle. Indeed, a
minimal increase around phase zero is only present for RV measurements based on iron in the period bin RRab1.

In light of these results, we opted to derive three independent mean RVCs for the Fe , the Mg , and the Na groups of lines. Selected RVCs for these three metallic diagnostics and for the four individual Balmer lines, anchored to $\tau_{0}=T_{\min }^{\mathrm{RV}(\mathrm{Fe})}$ and $\tau_{0}=T_{\min }^{\mathrm{RV}(\mathrm{H} \beta)}$, are shown in Figures 10 and 11 .

### 5.2. Analytical Fits of the RVC Templates

To construct the RVC templates, we normalized the RVCs by subtracting $V_{\gamma}$ and dividing by $\operatorname{Amp}\left(V_{r}\right)$. Once normalized, we stacked the RVCs within the same period bin of RVC template, thus obtaining the CNRVCs (see Appendix B).

To provide robust RVC templates, we decided to fit the CNRVCs with an analytical function. We discarded the Fourier series as a fitting curve because the number of phase points is not large enough (at least in the RRab3 period bin) to avoid


Figure 6. Left: Bailey diagram-optical amplitude vs. logarithmic period-for TS RRLs. Red squares and blue open circles display RRab and RRc variables. The vertical dashed lines separate the three period bins for RRab variables. Right: same as the left, but for iron $\operatorname{Amp}(V)$.


Figure 7. Same as Figure 2, but for the cumulative and normalized RVCs based on the Fe group lines. Panels (a), (c), (e), and (g) vs. panels (i), (k), (m), and (o) show the difference between RVCs phased by assuming as a reference epoch $\tau_{0}=T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and $\tau_{0}=T_{\text {min }}^{\mathrm{RV}(\mathrm{Fe})}$, respectively. Panels (b), (d), (f), (h), (j), (l), (n), and (p) display the residuals of the observations from the analytical fits.


Figure 8. Same as Figure 7, but for cumulative and normalized RVCs based on the $\mathrm{H}_{\beta}$ line. Note that panels (a)-(h) are anchored to $\tau_{0}=T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$ and panels (i)-(p) are anchored to $\tau_{0}=T_{\min }^{\mathrm{R} V(\mathrm{Fe})}$.
unphysical bumps and spurious secondary features in the fits. We adopted the PEGASUS fit as for the light-curve templates in Section 2.

Table 10 provides the coefficients of the PEGASUS functions obtained from the fitting procedure. When possible, we favored the lowest possible order, especially for the RRc template, as first-overtone pulsators have more sinusoidal RVCs. These are also the coefficients for the analytical form of the RVC templates. Figures 16 and 17 display the analytical fits of the RVC templates together with the observed RV measurements.

The largest standard deviations are those for the $\mathrm{H}_{\gamma}$ and $\mathrm{H}_{\delta}$ RVC templates for the RRc period bin $(\sim 0.08$ and 0.13 , respectively). To convert the $\sigma$ into the uncertainty on $V_{\gamma}$, one has simply to factor in $\operatorname{Amp}(\mathrm{RV})$. Since the typical $\operatorname{Amp}(\mathrm{RV})$ for RRc variables in $\mathrm{H}_{\gamma}$ and $\mathrm{H}_{\delta}$ range between 10 and $30 \mathrm{~km} \mathrm{~s}^{-1}$ (Bono et al. 2020b), the largest possible uncertainty introduced by the RVC template is of $\sim 4 \mathrm{~km} \mathrm{~s}^{-1}$. However, the largest absolute uncertainties are associated with the $\mathrm{H}_{\alpha}$ and $\mathrm{H}_{\beta}$ RVC templates for RRab1 period bin ( $\sigma \sim 0.05$ ). Since Amp (RV) is much larger for these diagnostics, the absolute uncertainties on $V_{\gamma}$ based on these RVC templates are of $\sim 7 \mathrm{~km} \mathrm{~s}^{-1}$. For metallic lines, all these uncertainties are on
average smaller than $\sim 3 \mathrm{~km} \mathrm{~s}^{-1}$. These results concerning both metallic and Balmer lines indicate that the current RVC templates can provide $V_{\gamma}$ for typical Halo RRLs with an accuracy better than $1 \%-3 \%$.

## 6. Reference Epoch to Apply the RVC Templates

A crucial aspect of templates is that they are used especially when the number of RV measurements is small. A first consequence is that, in a realistic scenario, the RV data is insufficient to accurately estimate $T_{\text {mean }}^{\mathrm{RV}}$ for either metal or Balmer RVCs. However, optical photometry usually is conducted before spectroscopic observations, and a good knowledge of the pulsation period and of $\operatorname{Amp}(V)$ are required to apply the RVC template. This means that we are typically dealing with a fairly well-sampled $V$-band light curve, and in turn, with an accurate estimate of $T_{\text {mean }}^{\text {opt }}$. With this in mind, it is necessary to check whether $T_{\text {mean }}^{\mathrm{RV}}$ and $T_{\text {mean }}^{\mathrm{opt}}$ in RRLs take place, within the errors, at the same phase along the pulsation cycle. If this is the case, we could safely make use of $T_{\text {mean }}^{\text {opt }}$ to phase the spectroscopic data and to apply the RVC template.
We phased a subset of RRLs with a good sampling of the pulsation cycle both in the $V$-band and in metallic RVCs. Fortunately, among the RRLs of the TS, there are several


Figure 9. Left: residuals of single-line Fe group RV measurements with respect to the mean RVC. Black (Fe1-mean), red (Fe2-mean), blue (Fe3-mean), and magenta (Sr-mean) symbols display the differences for the individual lines. The means of the residuals are displayed as solid lines of the same color of the symbols and are labeled together with their standard deviations. From top to bottom, residuals are displayed for the four period bins. Middle: same as the left, but for Mg RV measurements. Right: same as the left, but for Na RV measurements.
objects that were used for the BW analysis and for which there are available $V$-band light curves and RVCs collected at relatively close epochs (at most within $\sim 3 \mathrm{yr}$ ). This is an important advantage for this consistency test because possible changes in phases (phase drifts) and the effect of period
derivatives are small. We derived both $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and $T_{\text {mean }}^{\mathrm{opt}}$ and they are listed in Table 11.

Data listed in this table clearly show that $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and $T_{\text {mean }}^{\mathrm{opt}}$ trace, within the errors, the same phase along the pulsation cycle. Indeed, the average phase difference is $0.007 \pm 0.019$


Figure 10. First row: from left to right, radial velocity curves of the RRc variable CS Eri for the average $\mathrm{Fe}, \mathrm{Mg}$, and Na groups of lines. Second row: same as the top, but for the RRab1 variable HH Pup. Third row: same as the top, but for the RRab2 variable RV Oct. Fourth row: same as the top, but for the RRab3 variable AT Ser.
and always smaller than 0.05 (see column 6 in Table 11). ${ }^{24}$ This means the two reference epochs provide the same phasing. As a consequence, the photometric $T_{\text {mean }}^{\mathrm{opt}}$ can be safely adopted to anchor the RVC templates.

We already mentioned that there is a difference in phase between $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and $T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$. This means that, when adopting $T_{\text {mean }}^{\text {opt }}$ to use the template on Balmer lines, it is necessary to first shift the phases by an offset $\Delta \tau_{\mathrm{H} \beta}^{\mathrm{Fe}}=\Phi\left(T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}-T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}\right)$. For

[^1]this reason, we adopted the data listed in Table 6 and found a linear trend of $\Delta \tau_{\mathrm{H} \beta}^{\mathrm{Fe}}$ as a function of the pulsation period (see Figure 12). The plausibility of the phase difference between metallic and Balmer lines is further supported by the empirical evidence that the standard deviation of the relation is vanishing (0.008). Indeed, it is almost one order of magnitude smaller than the standard deviation of the phase offset between $T_{\text {min }}^{\mathrm{RV}(\mathrm{Fe})}$ and $T_{\min }^{\mathrm{RV}(\mathrm{H} \beta)}$ (see Section 4.5).

Large photometric surveys, however, often provide $T_{\text {max }}^{\text {opt }}$ but not $T_{\text {mean }}^{\text {opt }}$. To overcome this limitation and to facilitate the use of the RVC templates, we provide relations for $\Delta \Phi$ in Section 2.2 that allow the template user to easily convert the phases anchored on $\tau_{0}=T_{\max }^{\mathrm{opt}}$ into phases anchored on $\tau_{0}=T_{\text {mean }}^{\text {opt }}$.


Figure 11. First row: from left to right, radial velocity curves of the RRc variable CS Eri for individual Balmer lines $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\left.\mathrm{H}_{\delta}\right)$. Second row: same as the top, but for the RRab1 variable HH Pup. Third row: same as the top, but for the RRab2 variable RV Oct. Fourth row: same as the top, but for the RRab3 variable AT Ser.

## 7. Validation of the RVC Templates

A solid validation of the RVC templates requires that photometric and RV measurements be as close as possible in time. This methodological approach provides the unique opportunity to derive accurate photometric (pulsation period, $V$-band mean magnitude, $\operatorname{Amp}(V)$, reference epoch) and spectroscopic $\left(V_{\gamma}\right)$ properties of the RRLs adopted for the validation. We have already mentioned in Section 4.5 that the temporal proximity of both photometric and spectroscopic data is only available for a small number of field RRLs. To overcome this limitation, we decided to select from the calibrating sample one RRL per period bin: YZ Cap for the RRc, V Ind for the RRab1, W Crt for the RRab2, and AT Ser for the RRab3. We label these four RRLs as the Template

Validation Sample (TVS), and they are the benchmark for the RVC template validation.
Ideally, the validation should be done with an RRL sample independent from the one adopted to construct the RVC templates. However, we have verified that, by removing the four TVS RRLs from the calibrating RRLs, the coefficients of the analytical fits are only minimally affected. Note that the selection of the validating variable might bias the uncertainties of the result, because there is a small degree of internal variation in the shape of the RVCs of the different period bins. To investigate on a quantitative basis the dependence on the validating variable, we performed several tests by using each of the RRLs included in each period bin as the validating variable, one after the other. Interestingly enough, we found that the

Table 10
Coefficients and Standard Deviations of the PEGASUS Fits to the CNRVCs


(This table is available in machine-readable form.)

Table 11
Difference in the Reference Epoch between Light and Radial Velocity Curves

| Name | Type | Period <br> (days) | $\begin{aligned} & \hline \hline T_{\text {mean }}^{\mathrm{R} V(\mathrm{Fe})} \mathrm{HJD} \end{aligned}$ | $T_{\text {mean }}^{\text {opt }}$ | $\Delta \Phi$ | Ref. ${ }^{\text {a }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| DH Peg | RRc | 0.25551624 | 46667.5938 | 46684.6796 | 0.0239 | 0 |
| TV Boo | RRc | 0.31256 | 47220.7066 | 47227.5768 | 0.0260 | 1 |
| T Sex | RRc | 0.32468493 | 47129.6898 | 47226.4140 | 0.1229 | 1 |
| RS Boo | RRab | 0.37736549 | 46985.6766 | 46949.4482 | 0.0109 | 2 |
| AV Peg | RRab | 0.3903809 | 47130.3415 | 47123.3206 | -0.0427 | 1 |
| V0445 Oph | RRab | 0.397026 | 45842.2860 | 46980.9462 | -0.0111 | 3 |
| RR Gem | RRab | 0.3973 | 47220.6576 | 47226.5974 | 0.0498 | 1 |
| TW Her | RRab | 0.39959577 | 46925.8161 | 46947.3878 | 0.0168 | 2 |
| AR Per | RRab | 0.42556048 | 47128.3468 | 47123.6650 | 0.0019 | 1 |
| $V$ Ind | RRab | 0.47959915 | 57619.5115 | 47814.0720 | -0.0061 | 4 |
| BB Pup | RRab | 0.48055043 | 46136.1969 | 47192.9099 | 0.0052 | $5+6$ |
| TU UMa | RRab | 0.5569 | 47129.8698 | 47227.4521 | 0.0088 | 1 |
| SW Dra | RRab | 0.56967009 | 46519.6350 | 46496.2706 | 0.0140 | 7 |
| WY Ant | RRab | 0.57434364 | 46135.3591 | 47193.2805 | -0.0091 | $5+6$ |
| RX Eri | RRab | 0.58725159 | 47130.7581 | 47226.4736 | 0.0096 | 1 |
| RV Phe | RRab | 0.59641862 | 47054.1546 | 46305.0578 | -0.0075 | 8 |
| TT Lyn | RRab | 0.59744301 | 47129.6929 | 47123.1210 | 0.0003 | 1 |
| UU Cet | RRab | 0.60610163 | 47055.3286 | 47055.9229 | 0.0195 | 4 |
| W Tuc | RRab | 0.64224028 | 47057.8121 | 47493.2459 | 0.0046 | 4 |
| X Ari | RRab | 0.65117537 | 45640.8262 | 45639.5058 | 0.0278 | 9 |
| SU Dra | RRab | 0.66041178 | 47129.4805 | 47227.8778 | 0.0081 | 1 |
| VY Ser | RRab | 0.7141 | 44743.7769 | 47655.1533 | -0.0120 | 3 |

Notes. From left to right, the columns give the ID, the pulsation mode, the pulsation period, $T_{\text {mean }}^{\mathrm{RV}}, T_{\text {mean }}^{\mathrm{opt}}$, the difference in phase, and the reference. In column 7 , the references for the RVC and light curves adopted to derive $\Delta \Phi$ are listed in the following order:
${ }^{\text {a (0) Jones et al. 1988a; (1) Liu \& Janes 1989; (2) Jones et al. 1988b; (3) Fernley et al. 1990; (4) Clementini et al. 1990; (5 + 6) Skillen et al. 1993a, 1993b; (7) Jones }}$ et al. 1987a; (8) Cacciari et al. 1987; (9) Jones et al. 1987b.


Figure 12. $\Delta \tau_{\mathrm{H} \beta}^{\mathrm{Fe}}$ vs. period for the TS RRLs. The linear relation fitting the data is displayed as a solid line and labeled at the top, together with the standard deviation of the relation.
medians (see Tables 12 and 13) agree within $1 \sigma$, while the variation of the dispersions is on average smaller than $20 \%$. This means that the selection of the validating variable has a minimal impact on the accuracy of the validation.
We adopted the $V_{\gamma}$ of the TVS RRLs obtained by directly fitting their RVCs (see Tables 8 and 9) and assumed these as the best estimates for the systemic velocity $\left(V_{\gamma(\text { best })}\right)$ to be used in the validation process. However, to use the RVC template, we need to convert $\operatorname{Amp}(V)$ into $\operatorname{Amp}(\mathrm{RV})$ and then use the latter to rescale the normalized analytical function. The ratio between $\operatorname{Amp}(V)$ and $\operatorname{Amp}(\mathrm{RV})$ together with the equations for the conversion are thoroughly discussed in Appendix A.

The analytical form of the RVC templates can be used both as curves to be anchored to a single phase point and as
functions to fit a small number of phase points (three or more). Therefore, we followed two different paths for the validation process, based either on a single phase point approach (Section 7.1) or on multiple phase points (Section 7.2). The key idea is to estimate the accuracy of the light-curve templates from the difference $\Delta V_{\gamma}$ between $V_{\gamma(\text { best })}$ and the systemic velocities estimated by adopting the RVC template ( $V_{\gamma(\text { templ })}$ ). Furthermore, the accuracy of the current RVC templates is quantitatively compared, using TVS RRLs, with similar RVC templates available in the literature (S12).

### 7.1. Single Phase Point

As a first step, we generated a grid of 100 phase points for the seven RVCs $\left(\mathrm{Fe}, \mathrm{Mg}, \mathrm{Na}, \mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\mathrm{H}_{\delta}$, which we label, respectively, with a $j$ index running from 1 to 7). We interpolated the analytical fits of the RVCs for the TVS RRLs on an evenly spaced grid of phases $\left(\phi_{i}=[0.00,0.01, \ldots 0.99]\right.$, where $i$ runs from 1 to 100). For each $\phi_{i}$ and each $\operatorname{RVC}(j)$, we generated an RV measurement with the sum $\mathrm{RV}_{i j}=\mathrm{RV}\left(\operatorname{fit}\left(\phi_{i}\right)\right)_{j}+r \sigma$. The two addenda are: (i) $\operatorname{RV}\left(\operatorname{fit}\left(\phi_{i}\right)\right)_{j}$, which is the value of the fit of the $\operatorname{RVC}(j)$ interpolated at the phase $\phi_{i}$; and (ii) $r \sigma$, which simulates random noise, where $\sigma$ is the standard deviation of the phase points around the fit and $r$ is a random number extracted from a standard normal distribution.

We applied the RVC templates on each of the simulated points, thus deriving 100 estimates of the systemic velocity $\left(V_{\gamma(i, j)}^{\text {templ }}\right.$ ) for each $\operatorname{RVC}(j)$. To provide a quantitative comparison, these estimates were performed using both our own and the S12 RVC templates. The S12 RVC templates were not applied to YZ Cap, since they are valid only for RRab variables. In discussing the difference between our own RVC


Figure 13. Left: black crosses show the Fe RV measurements that we generated from the grid for the RRc variable YZ Cap. Gray dashed lines display the RVC templates from this work (labeled as TW) associated with each phase point. The ID of the RRL is labeled. Note that the comparison with the RVC template provided by S12 is missing because YZ Cap is an RRc variable. Middle: same as the left, but for the RRab variable $V$ Ind belonging to the RRab1 period bin. Top: radial velocity curves based on the RVC templates provided by S12. Bottom: same as the top, but based on our RVC templates. Right: same as the middle, but for the RRab variable SX For belonging to the RRab2 period bin.

Table 12
Validation of the RVC Templates Based on the Single Phase Point Approach

| Name | $\Delta V_{\gamma}(\mathrm{Fe})$ |  | $\Delta V_{\gamma}(\mathrm{Mg})$ |  | $\Delta V_{\gamma}(\mathrm{Na})$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\alpha}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\beta}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\gamma}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\delta}\right)$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{ }{ }^{\sigma}$ |  | $\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{ }{ }^{\sigma}$ |  | ${\underset{\left(\mathrm{km} \mathrm{~s}^{-1}\right)}{\mathrm{mdn}}}^{\sigma}$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  |
| Our RVC templates |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| YZ Cap | -0.082 | 2.270 | -0.548 | 2.755 | -0.970 | 2.600 | -0.815 | 3.027 | -1.116 | 3.131 | 0.202 | 2.566 | 0.062 | 2.906 |
| V Ind | -0.301 | 1.925 | -0.173 | 2.271 | 0.469 | 1.729 | -0.940 | 2.847 | 0.055 | 2.998 | 0.203 | 2.357 | -0.194 | 2.505 |
| SX For | 0.066 | 2.138 | -0.257 | 2.670 | -0.241 | 1.887 | 1.063 | 3.181 | 0.527 | 3.110 | 0.725 | 4.293 | 1.007 | 3.347 |
| S12 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| $V$ Ind | 0.049 | 2.977 | -0.583 | 3.287 | 0.440 | 3.200 | 0.685 | 4.356 | 0.877 | 6.601 | 0.644 | 6.895 | $\ldots$ | $\ldots$ |
| SX For | 1.488 | 5.418 | 1.115 | 5.583 | 1.467 | 5.634 | 3.863 | 8.812 | 3.820 | 10.926 | 4.154 | 10.899 | $\cdots$ | $\ldots$ |

Note. Medians (mdn) and standard deviations $(\sigma)$ for the $\Delta V_{\gamma}$ based on a single phase point validation, for both our and S12 RVC templates.

Table 13
Results of the Validation for the Multiple Phase Point Approach

| Name | $\Delta V_{\gamma}(\mathrm{Fe})$ |  | $\Delta V_{\gamma}(\mathrm{Mg})$ |  | $\Delta V_{\gamma}(\mathrm{Na})$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\alpha}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\beta}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\gamma}\right)$ |  | $\Delta V_{\gamma}\left(\mathrm{H}_{\delta}\right)$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  | $\left(\mathrm{km} \mathrm{~s}^{-1}\right)$ |  |
| Our RVC templates |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| YZ Cap | 0.028 | 3.049 | 1.682 | 2.857 | 0.873 | 3.666 | 4.147 | 5.813 | 0.742 | 3.863 | -1.702 | 5.413 | 4.839 | 4.799 |
| $V$ Ind | -0.472 | 5.658 | 0.808 | 4.649 | -0.621 | 6.230 | 6.268 | 17.383 | 3.227 | 11.646 | 2.620 | 9.894 | 2.978 | 8.356 |
| SX For | 1.436 | 8.031 | 2.943 | 8.068 | 1.840 | 9.490 | 4.373 | 16.630 | 1.700 | 12.225 | 3.050 | 10.995 | 4.196 | 9.721 |
| AT Ser | 0.219 | 9.340 | 2.311 | 11.129 | 0.363 | 12.087 | -0.583 | 17.990 | -0.844 | 15.311 | -0.371 | 13.432 | 4.477 | 12.727 |
| S12 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| V Ind | -0.273 | 7.822 | 0.770 | 9.314 | -0.777 | 9.805 | 4.817 | 19.250 | 2.676 | 13.259 | 2.263 | 12.345 | $\ldots$ | $\ldots$ |
| SX For | 2.297 | 7.673 | 3.953 | 7.502 | 2.437 | 9.102 | 3.988 | 16.418 | 2.207 | 11.638 | 3.522 | 10.763 | $\ldots$ | $\ldots$ |
| AT Ser | 0.454 | 9.522 | 2.412 | 10.869 | 0.528 | 11.037 | 1.251 | 19.826 | 0.989 | 16.109 | 1.289 | 13.684 | $\ldots$ | $\ldots$ |

Note. Medians (mdn) and standard deviations ( $\sigma$ ) for the $\Delta V_{\gamma}$ obtained from the three-point validation, for both our and the S12 RVC templates.
templates and those provided by S12, there are three key points worth being mentioned: (i) the comparison for $\mathrm{Fe}, \mathrm{Mg}$, and Na RVCs was performed by using the generic metallic RVC template from S12, because they do not provide elementspecific RVC templates; (ii) the RVC templates by S 12 were anchored to $T_{\max }^{\mathrm{opt}}$, while our own were anchored to $T_{\text {mean }}^{\mathrm{opt}}$; (iii) the RVC templates by S12 were rescaled, for internal consistency, by using their conversion equations from Amp $(V)$ to $A m p(R V)$ and not our own equation for the amplitude ratio.

Note that, due to the paucity of RRab3 variables, it was not possible to find one with both a reliable estimate of the reference epoch and with RV measurements close in epoch to the optical light curve. Therefore, the RVC templates for the RRab3 period bin were not validated with the single phase point approach. Nonetheless, we anticipate that we successfully validated the RRab3 RVC templates with the three-point approach (see Section 7.2). Figure 13 displays the simulated $\mathrm{RV}(\mathrm{Fe})$ phase points and the RVC templates applied to them. Finally, we derived the offsets $\Delta V_{\gamma(i, j)}=V_{\gamma(i, j)}^{\text {templ }}-V_{\gamma(j)}^{\text {best }}$ for each point and each RVC template. Table 12 gives the median and standard deviations of $\Delta V_{\gamma}$ for each RVC template.

We note that that the median of the $\Delta V$ is always smaller, in absolute value, than $1.0 \mathrm{~km} \mathrm{~s}^{-1}$ and $1.5 \mathrm{~km} \mathrm{~s}^{-1}$ for metallic and Balmer RVC templates. In all cases, the standard deviations are larger than the residuals, meaning that the latter can be considered vanishing within the dispersion. The largest standard deviations are found for the $\mathrm{H}_{\alpha}$ and $\mathrm{H}_{\beta}$ RVC templates and progressively decrease when moving to $\mathrm{H}_{\gamma}$, $\mathrm{H}_{\delta}$, and metallic lines.

The comparison between the current and the RVC templates provided by S12 makes it evident that the standard deviations of the former ones are systematically smaller by a factor ranging from $\sim 1.5$ to $\sim 3$. The higher accuracy of the current RVC templates is due to an interplay of a more robust estimate of $T_{\text {mean }}$ with respect to $T_{\max }$ and a more accurate optical-to-RV amplitude conversion (note, e.g., in the upper right panel of Figure 13, $\operatorname{Amp}(\mathrm{RV})$ is clearly overestimated for the S12 RVC template).

### 7.2. Multiple Phase Points

To apply the RVC templates to single phase points, it is necessary to know four parameters with great accuracy: (i) the pulsation period, (ii) the pulsation mode, (iii) $\operatorname{Amp}(\mathrm{V})$, and (iv) the reference epoch of the anchor point ( $T_{\text {mean }}^{\mathrm{opt}}$ ). The last one is particularly delicate, not only because a good sampling of the light curve is needed, but also because, when the spectroscopic data were collected several years before/after the photometric data, even small phase shifts or period changes can affect the phasing of the RV measurements. Note that, for RRLs affected by large Blazhko modulations, these parameters-especially $\operatorname{Amp}(V)$ and reference epoch-cannot be accurately estimated. Therefore, we suggest to avoid the application of the templates to RRLs with evident Blazhko modulations.

To overcome this limitation, it is possible to use the RVC templates as fitting functions if at least three RV measurements are available. In this empirical framework, only three parameters are required in order to apply the RVC template, namely the pulsation period, the pulsation mode, and $\operatorname{Amp}(V)$. We performed a number of tests by assuming that three independent RV measurements were available. In this
approach, the $\operatorname{Amp}(V)$ has to be converted into the Amp ( RV ) and then a least-squares fit of the RV measurements is performed by adopting the RVC template as the fitting function. The minimization of the $\chi^{2}$ is based on two parameters: $\Delta \phi$ (a horizontal shift) and $\Delta V_{\gamma}$ (a vertical shift). The function to be minimized is

$$
\begin{align*}
& P\left(\phi ; \Delta \phi, \Delta V_{\gamma}\right)=\Delta V_{\gamma}+\operatorname{Amp}(\mathrm{RV}) \\
& \quad\left(A_{0}^{P}+\Sigma_{i} A_{i}^{P} \exp \left(-\sin \left(\frac{\pi\left(\phi-\phi_{i}-\Delta \phi\right)}{\sigma_{i}^{P}}\right)^{2}\right)\right) . \tag{2}
\end{align*}
$$

To simulate three RV measurements, we extracted three random phases over the pulsation cycle and generated three RV measurements following the same approach adopted in Section 7.1. To rely on a wide statistical sample, the process was repeated 5000 times to generate 5000 different triplets of RV measurements for each RVC template. Henceforth, we label each of the triplets with a subscript $k$.

We estimated the systemic velocity $V_{\gamma(k, j)}^{\text {templ }}$ by fitting each of the triplets with both S12 and our own analytical form of the RVC templates. Analogously to Section 7.1, we derived the offsets $\Delta V_{\gamma(k, j)}$ and their median and standard deviations, and they are listed in Table 13.

We note that the uncertainties are larger for the three-point approach with respect to the single point one discussed in the former section. This happens because, when the three points are randomly extracted, it may happen that two (or all) of them are too close in phase, and the fitting procedure provides less solid estimates. This is a realistic case in which one might not have control over the sampling of the spectroscopic data. Also, the very fact that we are fitting the shift in phase-and not anchoring the RV data to the true reference epoch of the variable-means that the phasing of the template is not fixed, and this naturally leads to larger uncertainties. We have verified that, if the reference epochs were available, we would obtain standard deviations that are $\sim 10 \%-30 \%$ smaller and medians that are up to $\sim 50 \%$ smaller than in the case of the singleepoch approach.

Note that we did not put any restriction on which RV measurements were chosen to form a triplet. More specifically, we did not remove triplets in which two phase points were very close in phase, thus mimicking a realistic situation where either only two RV independent measurements were obtained or where the points are really close in phase due to the sampling of the spectroscopic data. Although this scenario usually leads to flawed estimates of $V_{\gamma}$, they are a minority, with less than $\sim 10 \%$ of the triplets having points closer than 0.05 in phase. Moreover, it is not only the difference in phase that matters, but also the distribution along the pulsation cycle. Indeed, phase points close in phase located along the decreasing branch of RRab variables produce larger uncertainties when compared with other phase intervals. The decision to keep even these troublesome triplets in our tests allowed the computed errors to take into account all the possible drawbacks of real observations, including the scenario of spectroscopic surveys that scan the sky by taking multiple consecutive measurements.

Data listed in Table 13 show that, with the only exception being RRab3, for which the standard deviations are similar, our RVC templates provide smaller standard deviations of the $\Delta V$ compared to S 12 templates. This is true even for the three-point validation, where the median offsets are smaller than the

Table 14
Properties of the RRLs Observed by MUSE in NGC 3201
$\left.\begin{array}{llccccccc}\hline \hline \text { Name } & \text { Type } & \begin{array}{c}\text { Period } \\ \text { (days) }\end{array} & V & \begin{array}{c}\text { Amp }(V) \\ (\mathrm{mag})\end{array} & \begin{array}{c}T_{\text {mean }}^{\text {opt }} \\ \text { HJD-2,400,000 } \\ \text { (days) }\end{array} & \begin{array}{c}T_{\text {max }}^{\text {opt }}\end{array} & V_{\gamma}\left(\mathrm{H}_{\alpha}\right) \\ (\mathrm{km} \mathrm{s}\end{array}\right)$
(This table is available in machine-readable form.)
standard deviations. This means that the RVC template provides $V_{\gamma}$ estimates that are more accurate than the simple average of the three RV measurements.

## 8. RR Lyrae in NGC 3201 as a Test Case

The Galactic globular cluster (GGC) NGC 3201 is a nearby ( 4.7 kpc ; Vasiliev \& Baumgardt 2021) metal intermediate ([Fe/ $\mathrm{H}] \sim-1.4$; Mucciarelli et al. 2014; Magurno et al. 2018) stellar system hosting a sizeable number of RRLs (77 RRab, 8 RRc, and 1 candidate RRd; Piersimoni et al. 2002; Layden \& Sarajedini 2003; Arellano Ferro et al. 2014) This cluster has a very high RV: $494.5 \pm 0.4 \mathrm{~km} \mathrm{~s}^{-1}$, based on giant and subgiant stars; and $496.47 \pm 0.11 \mathrm{~km} \mathrm{~s}^{-1}$, based on turn-off, subgiant, and red giant stars (Ferraro et al. 2018; Wan et al. 2021, respectively). Note that Magurno et al. (2018) adopted the S12 RVC templates and applied them to eleven RRLs in NGC 3201, obtaining a cluster RV of $494 \pm 2 \pm 8 \mathrm{~km} \mathrm{~s}^{-1}$, which is within $1 \sigma$ of literature estimates. Interestingly, by taking advantage of accurate proper motion measurements based on Gaia DR2, Massari et al. (2019) suggested that NGC 3201 is likely an object accreted during either the Sequoia or the Gaia-Enceladus merger.

To investigate on a more quantitative basis the cluster properties, we used the MUSE@VLT spectra collected by Giesers et al. (2019) for seven cluster RRab variables (see their properties in Table 14). From these spectra, we measured $\mathrm{H}_{\alpha}$ and $\mathrm{H}_{\beta}$ RVs (see data listed in Table 15) and the RVCs are displayed in Figure 14.

Unfortunately, the optical light curves that we adopted from Arellano Ferro et al. (2014) do not fully cover the pulsation for all the cluster RRLs. More specifically, we could not derive a reliable estimate of the reference epochs for V5 and V100. Due to this, we are not able to apply the templates to these two stars. In principle, we could apply the three-point method, but we wanted to keep our $V_{\gamma}$ estimates as homogeneous as possible.

After deriving-with the PLOESS algorithm-the best estimate for the systemic velocity ( $V_{\gamma \text { (best) }}$, displayed in Table 14) by fitting the RVCs, we applied the RVC template by using the single phase point approach. Note that, in this case, we do not generate a grid of synthetic points; we simply extract the points one by one from the RVC. Table 16 shows the average $V_{\gamma \text { (templ) }}$ estimates for all the lines, together with the standard deviation and the median of the offsets. The results are similar to those found in Section 7.1, meaning that all the medians are smaller than the standard deviations, and on average, the standard deviations obtained by using our own

Table 15
Radial Velocity Measurements for RRLs in NGC 3201

| Name | Species | HJD <br> (days) | $R^{\mathrm{a}}$ <br> $\left(\mathrm{km} \mathrm{s}^{-1}\right)$ | $\mathrm{eRV}^{\mathrm{b}}$ |
| :--- | :--- | :---: | :---: | :---: |
| V5 | Halpha | 2456978.8484 | 487.562 | 4.023 |
| V5 | Halpha | 2456989.8680 | 513.958 | 3.719 |
| V5 | Halpha | 2457008.8324 | 547.947 | 4.861 |
| V5 | Halpha | 2457009.8025 | 530.079 | 2.620 |
| V5 | Halpha | 2457131.4746 | 472.353 | 3.208 |
| V5 | Halpha | 2457134.4825 | 470.448 | 2.996 |
| V5 | Halpha | 2457138.4761 | 473.754 | 3.604 |
| V5 | Halpha | 2457419.7608 | 462.131 | 3.225 |
| V5 | Halpha | 2457421.7614 | 458.032 | 4.901 |
| V5 | Halpha | 2457787.8737 | 456.275 | 3.388 |

Notes. Only ten lines are listed. The machine-readable version of this table is available online on the CDS.
${ }^{\text {a }}$ Velocity plus heliocentric velocity and diurnal velocity correction.
${ }^{\mathrm{b}}$ Uncertainty on the RV measurements. For the Balmer lines, it is the uncertainty from spectroscopic data reduction. For $\mathrm{Fe}, \mathrm{Mg}$, and Na , it is the standard deviation of the RV measurements from different lines.
(This table is available in its entirety in machine-readable form.)

RVC templates are smaller than those coming from the use of the S12 RVC templates.

Our final estimate for the velocity of the whole system, by using both $\mathrm{H}_{\alpha}$ and $\mathrm{H}_{\beta}$ is $496.89 \pm 8.37$ (error) $\pm 3.43$ (standard deviation) $\mathrm{km} \mathrm{s}^{-1}$ by using the RVC templates and $495.21 \pm 3.23$ (error) $\pm 4.32$ (standard deviation) $\mathrm{km} \mathrm{s}^{-1}$ by simply fitting the RVCs, which agree quite well, within $1 \sigma$, with similar estimates available in the literature.

## 9. Summary and Final Remarks

We provide accurate and homogeneous RVC templates for both RRab and RRc variables using for the first time spectroscopic diagnostics based on well-identified metallic ( $\mathrm{Fe}, \mathrm{Mg}$, and Na ) lines and on Balmer $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\left.\mathrm{H}_{\delta}\right)$ lines. In the following, we discuss the approach adopted to construct the RVC templates, and we summarize the most relevant results.
$V$-band light-curve templates-To demonstrate on a quantitative basis the difference between the reference epoch anchored to the maximum brightness and that anchored to the mean magnitude along the rising branch, we collected homogeneous $V$-band photometry for cluster ( $\omega$ Cen and M4) and field RRLs. We have grouped them into four period bins (the same that we have used for the RVC templates) and found


Figure 14. Red: $\mathrm{H}_{\alpha}$ RV measurements. Blue: $\mathrm{H}_{\beta}$ RV measurements. The phases in the figures were derived by adopting the period and the reference epoch ( $T_{\text {mean }}^{\text {opt }}$ ) displayed in Table 14. For V5 and V100, we have assumed an arbitrary reference epoch, since we could not derive an accurate estimate from the optical light curves.

Table 16
Systemic Velocities of the RRLs in NGC 3201 Based on the RVC Templates

| Name | $\mathrm{H}_{\alpha}$ |  |  | $\mathrm{H}_{\beta}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $V_{\gamma(\text { templ })}$ | $\begin{aligned} & \sigma V_{\gamma(\text { templ })} \\ & \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{aligned}$ | $m d n \Delta V$ | $V_{\gamma(\text { templ })}$ | $\begin{gathered} \left.\sigma V_{\gamma(\text { templ })}\right) \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $m d n \Delta \mathrm{~V}$ |
|  | Our RVC templates |  |  |  |  |  |
| V17 | 500.4 | 6.0 | -0.6 | 499.4 | 5.6 | -2.3 |
| V23 | 495.7 | 6.1 | -0.6 | 498.1 | 8.8 | 4.9 |
| V50 | 498.0 | 6.3 | -1.0 | 498.0 | 5.5 | -0.6 |
| V77 | 486.9 | 12.2 | 3.6 | 489.4 | 8.0 | 1.4 |
| V90 | 491.1 | 10.3 | 1.3 | 494.8 | 10.7 | 1.8 |
|  | S12 |  |  |  |  |  |
| V17 | 496.2 | 15.3 | -4.8 | 498.5 | 9.4 | -3.2 |
| V23 | 498.6 | 9.2 | 2.2 | 497.5 | 10.2 | 4.3 |
| V50 | 499.3 | 5.3 | 0.3 | 497.1 | 7.3 | -1.5 |
| V77 | 488.3 | 12.1 | 5.1 | 488.3 | 11.7 | 0.3 |
| V90 | 492.3 | 18.3 | 2.5 | 494.1 | 22.3 | 1.1 |

(This table is available in machine-readable form.)
that the anchoring to the epoch of the mean magnitude on the rising branch ( $T_{\text {mean }}^{\mathrm{opt}}$ ) provides smaller standard deviations on the light-curve templates than the anchoring to the maximum brightness $\left(T_{\max }^{\mathrm{opt}}\right)$. The decrease is on the order of $35 \%$ for the period bins with cuspy light curves (RRab1 and RRab2) and of $45 \%$ for those with flat-topped light curves. These findings strongly support the results obtained by Inno et al. (2015) and
by Braga et al. (2019) in using $T_{\text {mean }}^{\text {opt }}$ to phase the NIR light curves of both classical Cepheids and RRLs.

Spectroscopic catalog-In this work, we present the largest collection of RRL spectra ever compiled in the literature. Altogether, we collected 23,865 spectra for 7070 RRab and 3343 RRc variables. These measurements were secured using eleven different spectrographs, ranging from low (2000) to very
high spectral resolution $(115,000)$. To build the RVC templates, the most important data set is the du Pont (see Table 3). Spectroscopic observations at this telescope were specifically planned to cover the entire pulsation cycle of several bright RRLs $(\langle V\rangle \sim 10-11 \mathrm{mag}$ for the majority of them). We also collected RV measurements and $V$-band time series from the literature (the BW sample), which were crucial to investigate the reference epoch and the ratio between RV and optical amplitudes.

Amplitude ratio-To apply the RVC templates, it is necessary to have prior knowledge on the optical amplitude of the variable, to correctly rescale the RVC template itself and to optimally match the true RVC. We provide new equations for the ratio of $\operatorname{Amp}(\mathrm{RV})$ over $\operatorname{Amp}(V)$. Those available in the literature ( S 12 ), for RVCs based on metallic and Balmer $\left(\mathrm{H}_{\alpha}\right.$, $\mathrm{H}_{\beta}$, and $\mathrm{H}_{\gamma}$ ) lines, were constructed using six RRab variables covering a very narrow range in pulsation period (0.56-0.59 days). In this investigation, we provide different RVC templates for both RRc and RRab variables based on metallic and Balmer $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\left.\mathrm{H}_{\delta}\right)$ lines. Even more importantly, our relations are based on three dozen variables covering a wide range in pulsation periods ( $0.27-0.83$ days) and metal content $(-2.6 \lesssim[\mathrm{Fe} / \mathrm{H}] \lesssim 0.0)$.

Reference epoch-When applying the RVC template to single RV measurements, it is necessary to anchor the RVC template at the same epoch as the observations. The RVC templates are applied to RRLs with well-sampled optical light curves and a few spectroscopic measurements. Therefore, the only pragmatic possibility to phase the spectroscopic data is to derive a reference epoch from the light curve. By using the light curves and the RV curves of our BW sample, we demonstrated that $T_{\text {mean }}^{\mathrm{opt}}$ is, within $5 \%$ of the pulsation cycle, identical to the time of mean velocity on the decreasing branch of the Fe RVC ( $\left.T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}\right)$. This means that RVC templates based on $\mathrm{Fe}, \mathrm{Mg}$, and Na lines can be safely anchored to $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$, and this approach does not introduce any systematic error when $T_{\text {mean }}^{\text {opt }}$ is adopted to apply the RVC template. We also found that RVCs based on Balmer lines display a well-defined phase lag across the phases of mean RV. Therefore, we decided to anchor the current $\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}$, and $\mathrm{H}_{\delta}$ RVC templates to the time of mean velocity on the decreasing branch of the $\mathrm{H}_{\beta}$ RVC $\left(T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}\right)$, taken as representative of the Balmer lines. Additionally, we found a clear trend in the phase difference between $T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and $T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$ as a function of the pulsation period. The new analytical relation gives the phase difference required to use the Balmer RVC templates, because $T_{\text {mean }}^{\mathrm{opt}}$ does not match $T_{\text {mean }}^{\mathrm{RV}(\mathrm{H} \beta)}$.

To discuss the concerning pros and cons of the reference epochs anchored to $T_{\text {mean }}^{\mathrm{RV}}$ and to $T_{\text {min }}^{\mathrm{RV}}$ in a more quantitative framework, we performed a series of numerical experiments. Interestingly enough, we found that RVC templates based on metallic RV measurements and anchored to $\tau_{0}=T_{\text {mean }}^{\mathrm{RV}}$ provide $V_{\gamma}$ that are on average $\sim 30 \%$ more accurate than the same RVCs anchored to $\tau_{0}=T_{\min }^{\mathrm{RV}}$. Even more importantly, we found that the improvement in using $T_{\text {mean }}^{\mathrm{RV}}$ compared with $T_{\text {min }}^{\mathrm{RV}}$ becomes even more relevant in dealing with RVCs based on Balmer lines. Indeed, the uncertainties decrease by up to a factor of three (see Section 7). The current circumstantial evidence further supports not only the use of two independent reference epochs for metallic and Balmer lines, but also the use of $T_{\text {mean }}^{\mathrm{RV}}$ as the optimal reference epoch to construct RVC templates.

Finally, we investigated the correlation in phase between $T_{\text {mean }}^{\mathrm{opt}}\left(\Phi_{\text {mean }}\right)$ and $T_{\text {max }}^{\mathrm{opt}}\left(\Phi_{\max }\right)$. This search was motivated by the fact that the current photometric surveys only provide reference epochs anchored to the time of maximum light $\left(T_{\max }^{\mathrm{opt}}\right)$. We found that the difference between the two $\left(\Delta \Phi_{\max -\min }=\Phi_{\max }-\Phi_{\min }\right)$ is constant-within the disper-sion-for RRc variables, and it follows a linear trend for RRab variables. We provide these relations for those interested in applying the current RVC templates to RRLs for which only $T_{\max }^{\mathrm{opt}}$ reference epochs are available.

RVC templates-We provide a total of 28 RVC templates of RRLs: these are divided into four different period bins (one for the RRc and three for the RRab variables) and seven diagnostics ( $\mathrm{Fe}, \mathrm{Mg}, \mathrm{Na}, \mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}$, and $\mathrm{H}_{\delta}$ ). The analytical form of the templates is provided in the form of a PEGASUS series (fifth to ninth order). Analytical fits based on the PEGASUS series, when compared with the Fourier series, do not display spurious ripples when there are no gaps in the phase coverage nor in the case of uneven sampling. The RVC templates have intrinsic dispersions that lead to errors smaller than $10 \mathrm{~km} \mathrm{~s}^{-1}$ in the worst case $\left(\mathrm{H}_{\alpha}\right.$ and $\mathrm{H}_{\beta}$ for highamplitude RRLs) and one order of magnitude smaller for the RVC templates with the smaller intrinsic dispersion (metallic lines RVC templates). To maximize the reach of the results of this work, we provide, in Appendix C, clear instructions on how to apply the RVC templates in different real-life observational scenarios.

Template validation-To validate the current RVC templates, we performed a detailed comparison with the RVC templates provided by S12 for RRab variables and based on metallic and Balmer $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}\right.$, and $\left.\mathrm{H}_{\gamma}\right)$ RVCs. We performed these tests on a subsample of variables that were used to build the RVC template (YZ Cap, V Ind, W Tuc, and AT Ser). The validation process was performed using both a single phase point approach, where the knowledge of the reference epoch is mandatory, and with a three-point approach, where the RVC template is used as a fitting function. We found that the median offset of the $V_{\gamma}$ estimates from the RVC templates is always smaller than $1.5 \mathrm{~km} \mathrm{~s}^{-1}$ (one-point approach) and $6 \mathrm{~km} \mathrm{~s}^{-1}$ (three-point approach). The medians are smaller than the standard deviations, meaning that systematic errors are negligible with respect to the statistical errors. We also found that our RVC templates provide $V_{\gamma}$ estimates that have a dispersion smaller by a factor of $1.5-3$ than those based on the RVC templates provided by S 12 .

RRLs in NGC 3201-We reduced the MUSE spectra already presented in Giesers et al. (2019), and obtained $\mathrm{H}_{\alpha}$ and $\mathrm{H}_{\beta}$ RVCs. We derived $V_{\gamma}$ both by fitting the RVC and extracting the measurements one by one, as well as by adopting the RVC templates. Their results based on these RV measurements are very similar to those of the validation process, with offsets smaller than $6 \mathrm{~km} \mathrm{~s}^{-1}$ and standard deviations that are smaller than those on the S12 RVC templates. Our estimate of the $V_{\gamma}$ of the cluster is $496.9 \pm 8.4$ (error) $\pm 3.4$ (standard deviation) $\mathrm{km} \mathrm{s}^{-1}$ from the RVC templates and $495.2 \pm 3.2$ (error) $\pm 4.3$ (standard deviation) $\mathrm{km} \mathrm{s}^{-1}$ from the fit with the RVC templates. They both agree, within $1 \sigma$, with literature estimates.

In the coming years, the current (OGLE, ASAS-SN, ZTF, sky-mapper, and PanSTARRS) and near-future (VRO) groundbased and space-based (Gaia, WISE, and WFIRST) observing facilities will provide a complete census of evolved variable stars associated with old stellar tracers in the Milky Way and in

Local Group galaxies. This will open new paths in the analysis of the early formation and evolution of the Galactic spheroid. However, firm constraints on the formation mechanism, in particular the dissipative collapse (Eggen 1962), dissipationless mechanism (Searle \& Zinn 1978), and Cold Dark Matter cosmological models (Monachesi et al. 2019), require detailed information concerning the kinematics and the metallicity distribution of the adopted stellar tracers. This is the approach already adopted to fully characterize the stellar streams and the merging history of the Galactic Halo (Helmi et al. 2018; Vasiliev 2019; Prudil et al. 2021).
Upcoming and ongoing low- (LAMOST-LR (Su et al. 1998), SDSS (Aguado et al. 2019)), medium- (4MOST (de Jong et al. 2012), SEGUE (Yanny et al. 2009), GALAH (Buder et al. 2018), HERBS (Duong et al. 2019), LAMOST-MR (Su et al. 1998), RAVE (Steinmetz et al. 2006), and WEAVE (Dalton et al. 2012)) and high-resolution (APOGEE (Majewski et al. 2017), MOONS (Taylor et al. 2018), and PFS (Tamura et al. 2018)) spectroscopic surveys will provide a wealth of new data for large samples of dwarf and giant field stars. In this context, old variable stars play a crucial role because their individual distances can be estimated with a precision on the order of $3 \%-$ $5 \%$ within the Local Group. Recent spectroscopic investigations based on high-resolution spectroscopy (For et al. 2011; Sneden et al. 2017; Crestani et al. 2021a) indicate that detailed abundance analysis can be performed with spectra collected at random phases. Unfortunately, the typical limiting magnitudes for spectroscopic investigations are roughly five magnitudes brighter than photometric ones, with current spectrographs available at the $8-10 \mathrm{~m}$ class telescope allowing us to reach limiting magnitudes on the order of $V \sim 20-21$ mag. However, even with the usage of large telescopes, the estimate of $V_{\gamma}$ is time-consuming, because it requires a spectroscopic time series of at least a dozen points. The RVC templates developed in this investigation provide new solid diagnostics to provide accurate $V_{\gamma}$ determinations by using a small number (three or less) of RV measurements, based on low-resolution spectra. Highly accurate estimates of line-of-sight velocities of stream stars are imperative for constraining the dark matter distribution (Bonaca \& Hogg 2018); these RVC templates will provide just that for the numerous RRL harbored by MW streams. The current diagnostics are focused on spectroscopic features located either in the blue or in the visual wavelength regime. A further extension into the red and the near-infrared regime is mandatory to fully exploit the most advanced space (Gaia (Gaia Collaboration et al. 2016), JWST (Gardner et al. 2006; Roman Spergel et al. 2015)) and ground-based (GMT, ELT, and TMT) observing facilities.

Based on observations made with ESO Telescopes at the La Silla Paranal Observatory under ESO programs 0100.D-0339, 0101.D-0697, 0102.D-0281, 076.B-0055, 077.B-0359, 077.D0633, 079.A-9015, 079.D-0262, 079.D-0462, 079.D-0567, 082.C-0617, 083.B-0281, 083.C-0244, 094.B-0409, 095.B0744, 097.A-9032, 098.D-0230, 189.B-0925, 267.C-5719, 297.D-5047, 67.D-0321, 67.D-0554, 69.C-0423, 71.C-0097, 0100.D-0273, 083.C-0244, 098.D-0230, 095.D-0629, 096.D0175, 097.D-0295, 098.D-0148, 0100.D-0161, 0101.D-0268, 0102.D0270, and 0103.D-0204.

Based on observations made with the Italian Telescopio Nazionale Galileo (TNG) operated on the island of La Palma by the Fundación Galieo Galilei of the INAF (Istituto

Nazionale di Astrofisica) at the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofisica de Canarias.
Some of the observations reported in this paper were obtained with the Southern African Large Telescope (SALT).
Funding for the SDSS and SDSS-II has been provided by the Alfred P. Sloan Foundation, the Participating Institutions, the National Science Foundation, the U.S. Department of Energy, the National Aeronautics and Space Administration, the Japanese Monbukagakusho, the Max Planck Society, and the Higher Education Funding Council for England. The SDSS Web Site is http://www.sdss.org/. The SDSS is managed by the Astrophysical Research Consortium for the Participating Institutions. The Participating Institutions are the American Museum of Natural History, Astrophysical Institute Potsdam, University of Basel, University of Cambridge, Case Western Reserve University, University of Chicago, Drexel University, Fermilab, the Institute for Advanced Study, the Japan Participation Group, Johns Hopkins University, the Joint Institute for Nuclear Astrophysics, the Kavli Institute for Particle Astrophysics and Cosmology, the Korean Scientist Group, the Chinese Academy of Sciences (LAMOST), Los Alamos National Laboratory, the Max-Planck-Institute for Astronomy (MPIA), the Max-Planck-Institute for Astrophysics (MPA), New Mexico State University, the Ohio State University, University of Pittsburgh, University of Portsmouth, Princeton University, the United States Naval Observatory, and the University of Washington.
V.F.B. acknowledges the financial support of the Istituto Nazionale di Astrofisica (INAF), Osservatorio Astronomico di Roma, and Agenzia Spaziale Italiana (ASI) under contract to INAF: ASI 2014-049-R. 0 dedicated to SSDC.
We acknowledge financial support from the US NSF under Grants AST-1714534 (M.M., J.P.M.) and AST1616040 (C.S.). B.L., Z.P., and H.L. were supported by the Deutsche Forschungsgemeinschaft (DFG; German Research Founda-tion)—Project-ID 138713538-SFB 881 ("The Milky Way System," subprojects A03, A05, and A11). E.V. acknowledges the Excellence Cluster ORIGINS Funded by the Deutsche Forschungsgemeinschaft (DFG; German Research Foundation) under Germany's Excellence Strategy-EXC-2094390783311. Z.P. gratefully acknowledges funding by the Deutsche Forschungsgemeinschaft (DFG; German Research Foundation)-Project-ID 138713538-SFB 881 ("The Milky Way System"), subprojects A03, A05, and A11. G.F. has been supported by the Futuro in Ricerca 2013 (grant RBFR13J716). M.L. acknowledges funding from the Deutsche Forschungsgemeinschaft (grant DR 281/35-1).
N.M. is grateful to the Japan Society for the Promotion of Science (JSPS) for the Grant-in-Aid for Scientific Research, KAKENHI 18H01248.

## Appendix A Amplitude Ratios

The fundamental piece of information that is needed to use the RVC template RVCs is the relation of $\operatorname{Amp}(\mathrm{RV})$ versus $\operatorname{Amp}(V)$. These relations should be adopted to rescale the normalized templates to the true $\operatorname{Amp}(\mathrm{RV})$ of the star, with the knowledge of $\operatorname{Amp}(V)$. The leading argument is always that, since there is far more availability of light curves than of RVCs, the amplitudes of the former should be adopted to derive those of the latter.

For their RVC templates, S12 adopted $\mathrm{Amp}(\mathrm{RV})$ versus $\operatorname{Amp}(V)$ relations based on six RRLs for metallic lines, $\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}$, and $\mathrm{H}_{\gamma}$. We provide the same relations, but separating RRc and RRab variables. We have considered the use of $\operatorname{Amp}(\mathrm{RV})$ versus $\operatorname{Amp}([3.6])$ relations, since the [3.6] band traces radius variations and is not affected by temperature variations as the $V$ band is. However, there are two arguments against this option: (i) there is a paucity of MIR light curves with respect to optical light curves; and (ii) there is mounting empirical evidence that amplitude ratios using NIR and MIR bands follow a mild trend with metallicity (Mullen et al. 2021).

Before deriving the relation between $\operatorname{Amp}(\mathrm{RV})$ and Amp $(V)$, we checked whether we could assume the $\mathrm{Fe}, \mathrm{Mg}$, and Na amplitudes as equivalent. Figure 15 shows the residuals of Na and $\operatorname{Mg} \operatorname{Amp}(\mathrm{RV})$ versus the $\mathrm{Fe} \operatorname{Amp}(\mathrm{RV})$. It is clear that there is no trend with period, and also the offset is well within the $\mathrm{Amp}(\mathrm{RV})$ uncertainty. This means that we can provide only two $\operatorname{Amp}(\mathrm{RV})$ versus $\operatorname{Amp}(V)$ relations (one for RRc and one for RRab variables) that hold for both $\mathrm{Fe}, \mathrm{Mg}$, and Na . We also collected RV curves of RRLs from literature, all derived from metallic lines (BW sample; see caption of Table 11). We found that the $\mathrm{BW} \operatorname{Amp}(\mathrm{RV})$ displays a very small difference with


Figure 15. Top left: difference between radial velocity amplitudes based on Mg and Fe group RVC templates. The mean difference is labeled at the top right corner. Middle left: same as top left, but for Na and Fe group RVC templates. Bottom left: same as the top left, but between the BW sample and the Fe group RVC templates. Right panels: from top to bottom, visual amplitude $(\operatorname{Amp}(V))$ vs. radial velocity amplitudes $(\mathrm{Amp}(\mathrm{RV}))$ based on $\mathrm{Fe}, \mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}$, and $\mathrm{H}_{\delta}$ lines. RRab and RRc variables are marked with red and blue symbols. Uncertainties on $\operatorname{Amp}(\mathrm{RV})$ are displayed. Solid lines show the linear relations between $\mathrm{Amp}(\mathrm{RV})$ and $\mathrm{Amp}(V)$. The analytical forms of the relations are labeled at the bottom.

Table 17
Coefficients of the $\operatorname{Amp}(\mathrm{RV})-\operatorname{Amp}(V)$ Relations

| Line | RRc |  |  |  |  | RRab |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{km} \mathrm{s}^{-1}$ |  | $\mathrm{km} \mathrm{s}^{-1} \mathrm{mag}^{-1}$ |  | $\sigma$ | $\mathrm{km} \mathrm{s}^{-1}$ |  | $\mathrm{km} \mathrm{s}^{-1} \mathrm{mag}^{-1}$ |  |  |
| Fe | 1.68 | 4.92 | 52.63 | 12.94 | 3.50 | 38.09 | 2.90 | 22.35 | 3.12 | 4.20 |
| $\mathrm{H} \alpha$ | -9.18 | 9.36 | 106.69 | 24.61 | 6.66 | 77.77 | 6.17 | 22.83 | 6.64 | 8.92 |
| $\mathrm{H} \beta$ | -0.27 | 6.92 | 66.55 | 18.20 | 4.93 | 63.32 | 4.96 | 15.96 | 5.34 | 7.17 |
| H $\gamma$ | 0.54 | 7.99 | 59.32 | 21.02 | 5.69 | 57.38 | 5.01 | 18.48 | 5.39 | 7.24 |
| $\mathrm{H} \delta$ | 3.85 | 6.80 | 45.78 | 17.88 | 4.84 | 50.90 | 3.84 | 14.43 | 4.14 | 5.55 |

Note. Coefficients $(a, b)$, uncertainties ( $e a, e b$ ), and standard deviations $(\sigma)$ of the $\operatorname{Amp}(\mathrm{RV})$ versus $\operatorname{Amp}(V)$ relations for RRc and RRab variables.
(This table is available in machine-readable form.)
our $\operatorname{Fe} \operatorname{Amp}(\mathrm{RV})\left(0.23 \pm 3.73 \mathrm{~km} \mathrm{~s}^{-1}\right)$. This allowed us to merge the two sets of $A m p(R V)$ and derive a more solid relation for the metallic $\operatorname{Amp}(\mathrm{RV})$, based on a larger number of objects ( 12 RRc and 60 RRab ).

The right panels of Figure 15 display the trend of Fe and of the Balmer $\operatorname{Amp}(\mathrm{RV})$ with $\operatorname{Amp}(V)$. A steady increase of the slopes for the Balmer $\operatorname{Amp}(\mathrm{RV})$, from $\delta$ to $\alpha$, is clear enough both visually in the figure and quantitatively from the coefficients listed in Table 17. This is also the first time that different trends have been found for RRc and RRab.

## Appendix B The Cumulated RVCs

Figures 16 and 17 display the CNRVCs for all the RVC templates (colored small circles), together with the analytical form of the template (solid line).

Inspecting the CNRVCs makes it clear that, within the bins RRab1 and RRab2, the morphology of the RVCs is, to a first approximation, dichotomic. More specifically, within the RRabl bin, SW Aqr, ST Leo, VX Her, V Ind, and V0440 Sgr display a local maximum around the phase $0.70-0.75$, instead of the more or less steady rising behavior of the other RRabl RRLs. Moreover, TY Gru, CD Vel, and SX For (RRab2 bin), do not display a local minimum around phase $0.7-0.8$, as
the other RRab2 variables do. This happens for all the diagnostics, although it is more evident for the Balmer lines. We remark that these features are also present in the optical light curves of these stars.

We checked whether these features can be associated with either pulsation or physical properties of the stars. While it is true that the RRabl variables with a more prominent local maximum are located in the High-Amplitude Short-Period (HASP; Fiorentino et al. 2015) region, they are not the only HASP variables in our sample. Their iron abundances range between -1.9 and -1.4 , which is around the peak of the distribution of field RRLs (Crestani et al. 2021a) and it was not even possible to constrain a morphological class of RVCs based on the Fourier parameters R2, R31, $\phi_{21}$, and $\phi_{31}$ of their light curve.

To sum up, there is no quantitative way to predict, either from the light curve or from the physical properties, which is the RVC morphology of RRab1 and RRab2. This has several consequences: the first and most obvious is that we cannot split these bins and provide more RVC templates, because we cannot provide criteria for using one or the other. This may seem to be a disadvantage, but luckily, this dichotomy introduces a $\lesssim 5 \mathrm{~km} \mathrm{~s}^{-1}$ offset in the estimate of $V_{\gamma}$ with a probability of $\sim 10 \%-15 \%$ (i.e., the fraction of pulsation cycle where the RVCs have a different behavior).




Figure 16. From top to bottom: cumulative and normalized radial velocity curves based on metallic ( $\mathrm{Fe}, \mathrm{Mg}$, and Na ) lines. The period bin of the RVC template is labeled on the top left corner. Small circles are color-coded by variables and their names are labeled at the bottom. The solid line displays the analytical form of the RVC templates.


Figure 17. From top to bottom: cumulative and normalized radial velocity curves based on the Balmer $\left(\mathrm{H}_{\alpha}, \mathrm{H}_{\beta}, \mathrm{H}_{\gamma}\right.$, and $\left.\mathrm{H}_{\delta}\right)$ lines. The period bin of the RVC template bin is labeled on the top left corner. Small circles are color-coded by variables and their names are labeled at the bottom. The solid line displays the analytical form of the RVC templates.

## Appendix C <br> How to Use the New RVC Templates

In this section of the appendices, we provide precise instructions on how to use the new RVC templates with the aim of estimating $V_{\gamma}$ in different realistic cases: (i) when only one RV measurement and a well-sampled light curve that allows the estimate of the reference epoch are available; (ii) when a few (we assume three) RV measurements and a wellsampled light curve that allows the estimate of the reference epoch are available; (iii) when a few (we assume three) RV measurements and a light curve that does not allow an accurate estimate of the reference epoch are available.

## C.1. Estimate of $T_{\text {mean }}^{\mathrm{opt}}$

Before describing how to use the templates, we want to instruct the reader on how to derive $T_{\text {mean }}^{\mathrm{opt}}$, which is not as common as $T_{\max }^{\mathrm{opt}}$, and therefore it might not be straightforward to estimate.

First of all, the light curve must be phased to an arbitrary reference epoch, e.g., HJD $=0$, as in Figure 18. Then, the phased light curve must be fit with a model. It is crucial that the model fits well the rising branch. After that, $\langle V\rangle$ is derived by converting each point of the model into arbitrary flux ( $F=10^{-0.4 * \mathrm{mag}}$ ), integrating the flux model, finding the mean flux $\langle F\rangle$, and converting back to magnitude $\left(\langle V\rangle=-2.5 \cdot \log _{10}(\langle F\rangle)\right)$.

If the model is analytical (e.g., Fourier or PEGASUS), one can easily find the phase on the rising branch at which the model intersects $\langle V\rangle$, namely $\phi_{\text {mean }}$. If the model is not analytical (e.g., spline or PLOESS), it is necessary to interpolate $\langle V\rangle$ with the model sampled on an even grid of phases. A convenient choice for the step of the grid is between 0.001 and 0.01 .

Once $\phi_{\text {mean }}$ is known, the next step is to select any phase point of the light curve. This will be characterized by an epoch of observation $\left(t_{V(i)}\right.$, where $i$ indicates the $i$ th point of the light curve) and a phase $\left(\phi_{V(i)}\right)$. Finally, $T_{\text {mean }}^{\text {opt }}=t_{V(i)}-\left(\phi_{V(i)}-\phi_{\text {mean }}\right) \cdot P$ can be derived.

## C.2. Single RV Measurement

This is the most classical situation when using any type of template. In this case, only one RV measurement is available. Note that, if the spectral range of the instrument is large enough, it is possible to have one RV measurement per diagnostic (e.g., $\mathrm{Fe}, \mathrm{Mg}$, and $\mathrm{H}_{\gamma}$ ), but still no more than one epoch is available. This means that any RVC template can be applied to only one point.
The quoted RV measurement consists of an epoch of observation ( $t$ ), a velocity (RV), and its uncertainty (eRV). In this case, a decent sampling of the light curve is needed in order to estimate its period $(P)$, amplitude $(\operatorname{Amp}(V))$, mean magnitude $(\langle V\rangle)$, and the reference epoch $T_{\max }^{\mathrm{opt}}$ or $T_{\text {mean }}^{\mathrm{opt}}$.

The first step is to anchor the RV measurement to the same reference epoch as the template. As demonstrated in the body of the paper, it is possible to assume $T_{\text {mean }}^{\mathrm{opt}}=T_{\text {mean }}^{\mathrm{RV}(\mathrm{Fe})}$ and derive the phase as $\phi=\frac{t-T_{\text {mean }}^{\mathrm{opt}}}{P} \bmod 1$. If only $T_{\max }^{\mathrm{opt}}$ estimates are available, as is often the case with data releases of large surveys, it is necessary to derive the phase anchored to $T_{\max }^{\mathrm{opt}}$ and then apply the offsets provided in Section 2.2: $\phi_{\max }=\frac{t-T_{\max }^{\mathrm{opt}}}{P} \bmod 1$ and then $\phi=\phi_{\max }+0.223$ for RRc


Figure 18. $V$-band ASAS light curve of SX For. The solid red line displays the fitting model. The dashed red line represents the mean magnitude, and the blue circle displays the position of $T_{\text {mean }}^{\text {opt }}$. The values of the mean magnitude and $T_{\text {mean }}^{\mathrm{opt}}$ are labeled in red and blue, respectively.
variables, or $\phi=\phi_{\max }+0.043+0.099 \cdot P$ for RRab variables. In this case, an uncertainty must be associated with $\phi$, namely the $\sigma$ of the quoted relations: 0.036 for RRc and 0.024 for RRab. By using $\phi$, the RV measurement is now anchored to the same reference of metallic ( $\mathrm{Fe}, \mathrm{Mg}$, and Na ) templates. If, however, the RV measurement is from a Balmer line, it is necessary to convert $\phi$ by using the relation provided in Figure 12, i.e., $\phi_{H \beta}=\phi+0.023-0.096 * P$.

The second step is to rescale of the normalized template. For this end, it suffices to convert $\operatorname{Amp}(V)$ into $\operatorname{Amp}(\mathrm{RV})$ by using the relations provided in Table 17.

The third step is to derive the analytical form of the template both rescaled by $\operatorname{Amp}(\mathrm{RV})$ and shifted in zero point to pass through the RV measurement. For this step, the right coefficients for the RVC template must be selected from Table 10. Afterward, these coefficients are substituted into Equation (1) to calculate the value of the template $(T(\phi))$ at the phase $\phi$. Finally, $V_{\gamma}$ is obtained as $V_{\gamma}=\mathrm{RV}-\operatorname{Amp}(\mathrm{RV}) \cdot T(\phi)$.

Of course, if RV measurements are available for more than one diagnostic, the quoted steps can be separately applied to all the diagnostics, providing multiple $V_{\gamma}$ estimates that can be averaged.

## C.3. Multiple RV Measurement with Reference Epoch

In this case, more than one RV measurement per diagnostic and a sufficiently well-sampled light curve are available. The procedure is qualitatively identical to that described in C.2, but having more than one RV measurement per diagnostic allows the averaging of the $V_{\gamma}$ estimates for the same diagnostic. Also in this case, if RV measurements are available for more than one diagnostic, these can be averaged to constrain $V_{\gamma}$ on a broader statistical basis.

In principle, this method could be applied to any number of RV measurements. However, when there are ten to twelve (or more) of these and they are more or less evenly sampled, with a good knowledge of the period, it is possible to just fit the points and directly derive a $V_{\gamma}$ estimate as accurate as the template itself-or even more so if the variable has experienced some phase drift or period change during the time elapsed from the collection of light curve and RV data.

## C.4. Multiple RV Measurement without Reference Epoch

In this case, more than two RV measurements per diagnostic are available, but the light curve is only modestly sampled. This means that period and $\operatorname{Amp}(V)$ can be estimated from photometric data, but the reference epoch cannot. In this situation, the templates can be used as fitting curves.

This approach is qualitatively different from the other two because one does not have to anchor the template but rather to fit it to the data. The steps to do so are as follows. First of all, phases must be derived by adopting an arbitrary reference epoch $T_{0}$ (e.g., $T_{0}=0$ or $\left.T_{0}=2,400,000\right) \phi=\frac{t-T_{0}}{P} \bmod 1$. Still, if one wants to use the Balmer templates, the conversion $\phi_{H \beta}=\phi+0.023-0.096 * P$ must be applied.
Second, $\operatorname{Amp}(V)$ has to be rescaled into $\operatorname{Amp}(\mathrm{RV})$ by using the relations provided in Table 17. This step is analogous to that described in C.2.

The third step is to select the RVC template coefficients from Table 10 and perform a $\chi^{2}$ minimization when fitting the RV data with Equation (2). The minimization must be performed on the two free parameters $\Delta \phi$ (a horizontal shift) and $\Delta V_{\gamma}$, while all the others remain fixed.

Finally, $V_{\gamma}$ is simply derived by integrating the final analytical form of Equation (2). Also, in this case, if RV measurements are available for more than one diagnostic, these can be averaged to constrain $V_{\gamma}$ on a broader statistical basis.

## Appendix D <br> Phase-gridded Templates

The anonymous referee suggested that the reader might be interested in using analytical functions different from the PEGASUS series to apply to RVC templates. To facilitate the

Table 18
Phase-gridded RVC Templates

| Species | Template Bin | Phase | Template Value $^{\mathrm{a}}$ |
| :--- | :---: | :---: | :---: |
| Fe | 0 | 0.000 | 0.0176 |
| Fe | 0 | 0.001 | 0.0124 |
| Fe | 0 | 0.002 | 0.0072 |
| Fe | 0 | 0.003 | 0.0019 |
| Fe | 0 | 0.004 | -0.0036 |
| Fe | 0 | 0.005 | -0.0092 |
| Fe | 0 | 0.006 | -0.0148 |
| Fe | 0 | 0.007 | -0.0206 |
| Fe | 0 | 0.008 | -0.0265 |
| Fe | 0 | 0.009 | -0.0325 |

Notes.
${ }^{\text {a }}$ Calculated at the given phase by using Equation (1) and the coefficients provided in Table 10.
(This table is available in its entirety in machine-readable form.)
independent readers, we list in Table 18 the phase-gridded values of the RVC templates by using the coefficients of the PEGASUS series given in Table 10. The reader can fit these gridded values with any analytical function and use them to apply the RVC template. We adopted a step of 0.01 ; this means 101 points per template. We have checked that, for a Fourier series of tenth-to-fifteenth order, this step in phase provides a very good fit of the gridded points. In passing, we stress that these phase-gridded values should not be used in substitution of the analytical templates, but to derive alternative analytical forms of the RVC templates. This is recommended for the single phase point approach and necessary for the multiple phase point approach.

## ORCID iDs

V. F. Braga (i) https://orcid.org/0000-0001-7511-2830
J. Crestani © https:// orcid.org/0000-0001-8926-3496
M. Fabrizio (i) https:// orcid.org/0000-0001-5829-111X
G. Bono (i) https:// orcid.org/0000-0002-4896-8841
C. Sneden (©) https:// orcid.org/0000-0002-3456-5929
J. Storm (10 https:// orcid.org/0000-0002-8627-6096
S. Kamann (10) https:// orcid.org/0000-0001-6604-0505
M. Latour (ib https://orcid.org/0000-0002-7547-6180
Z. Prudil © https:// orcid.org/0000-0001-5497-5805
G. Altavilla (i) https://orcid.org/0000-0002-9934-1352
B. Chaboyer © https://orcid.org/0000-0003-3096-4161
M. Dall'Ora © https:// orcid.org/0000-0001-8209-0449
I. Ferraro © https:// orcid.org/0000-0001-8514-7957
C. K. Gilligan (©) https:// orcid.org/0000-0003-4510-0964
G. Fiorentino (©) https:// orcid.org/0000-0003-0376-6928
G. Iannicola © https:// orcid.org/0000-0001-9816-5484
L. Inno © https://orcid.org/0000-0002-0271-2664
M. Marengo (1) https:// orcid.org/0000-0001-9910-9230
S. Marinoni © https://orcid.org/0000-0001-7990-6849
P. M. Marrese (1) https:// orcid.org/0000-0002-8162-3810
C. E. Martínez-Vázquez © https://orcid.org/0000-0002-9144-7726
M. Monelli © https://orcid.org/0000-0001-5292-6380
J. P. Mullen © https://orcid.org/0000-0002-1650-2764
J. Neeley © https:// orcid.org/0000-0002-8894-836X
P. B. Stetson (1) https://orcid.org/0000-0001-6074-6830
E. Valenti © https://orcid.org/0000-0002-6092-7145
M. Zoccali (1) https://orcid.org/0000-0002-5829-2267

## References

Aguado, D. S., Ahumada, R., Almeida, A., et al. 2019, ApJS, 240, 23 Arellano Ferro, A., Ahumada, J. A., Calderón, J. H., \& Kains, N. 2014, 50, 307
Baade, W. 1926, AN, 228, 359
Baade, W. 1956, PASP, 68, 5
Belokurov, V., Deason, A. J., Koposov, S. E., et al. 2018, MNRAS, 477, 1472
Benkő, J. M., Sódor, Á., \& Pál, A. 2021, MNRAS, 500, 2554
Bonaca, A., \& Hogg, D. W. 2018, ApJ, 867, 101
Bono, G., Caputo, F., Cassisi, S., Incerpi, R., \& Marconi, M. 1997, ApJ, 483, 811
Bono, G., Caputo, F., Castellani, V., et al. 2003, MNRAS, 344, 1097
Bono, G., Caputo, F., \& Stellingwerf, R. F. 1994, ApJL, 432, L51
Bono, G., \& Cassisi, S. 1999, in ASP Conf. Ser. 165, The Third Stromlo Symposium: The Galactic Halo, ed. B. K. Gibson, R. S. Axelrod, \&
M. E. Putman (San Francisco, CA: ASP), 290, arXiv:astro-ph/9811175

Bono, G., Braga, V. F., Crestani, J., et al. 2020a, ApJL, 896, L15
Bono, G., Braga, V. F., Fiorentino, G., et al. 2020b, A\&A, 644, A96
Braga, V. F., Bhardwaj, A., Contreras Ramos, R., et al. 2018, A\&A, 619, A51
Braga, V. F., Bono, G., Fiorentino, G., et al. 2020, A\&A, 644, A95
Braga, V. F., Dall'Ora, M., Bono, G., et al. 2015, ApJ, 799, 165
Braga, V. F., Stetson, P. B., Bono, G., et al. 2016, AJ, 152, 170
Braga, V. F., Stetson, P. B., Bono, G., et al. 2019, A\&A, 625, A1
Buder, S., Asplund, M., Duong, L., et al. 2018, MNRAS, 478, 4513
Cacciari, C., Clementini, G., Prevot, L., et al. 1987, Åps, 69, 135
Carney, B. W., Storm, J., \& Jones, R. V. 1992, ApJ, 386, 663
Carollo, D., Beers, T. C., Lee, Y. S., et al. 2007, Natur, 450, 1020
Castor, J. I. 1971, ApJ, 166, 109
Catelan, M. 2009, Ap\&SS, 320, 261
Catelan, M., Pritzl, B. J., \& Smith, H. A. 2004, ApJS, 154, 633
Clementini, G., Cacciari, C., \& Lindgren, H. 1990, A\&AS, 85, 865
Clementini, G., Ripepi, V., Molinaro, R., et al. 2019, A\&A, 622, A60
Conroy, C., Naidu, R. P., Garavito-Camargo, N., et al. 2021, arXiv:2104.09515
Crestani, J., Fabrizio, M., Braga, V. F., et al. 2021a, ApJ, 908, 20
Crestani, J., Braga, V. F., Fabrizio, M., et al. 2021b, arXiv:2104.08113
Dalton, G., Trager, S. C., Abrams, D. C., et al. 2012, Proc. SPIE, 8446, 220
de Jong, R. S., Bellido-Tirado, O., Chiappini, C., et al. 2012, Proc. SPIE, 8446, 252

Deason, A. J., Belokurov, V., \& Evans, N. W. 2011, MNRAS, 416, 2903
Drake, A. J., Djorgovski, S. G., Mahabal, A., et al. 2009, ApJ, 696, 870
Drake, A. J., Djorgovski, S. G., Catelan, M., et al. 2017, MNRAS, 469, 3688 Duong, L., Asplund, M., Nataf, D. M., et al. 2019, MNRAS, 486, 3586
Eggen, O. J. 1962, PASP, 74, 159
Fabrizio, M., Bono, G., Braga, V. F., et al. 2019, ApJ, 882, 169
Fernley, J. A., Skillen, I., Jameson, R. F., et al. 1990, MNRAS, 247, 287
Ferraro, F. R., Mucciarelli, A., Lanzoni, B., et al. 2018, ApJ, 860, 50
Fiorentino, G., Bono, G., Monelli, M., et al. 2015, ApJL, 798, L12
For, B.-Q., Sneden, C., \& Preston, G. W. 2011, ApJS, 197, 29
Gaia Collaboration, Prusti, T., de Bruijne, J. H. J., et al. 2016, A\&A, 595, A1
Gardner, J. P., Mather, J. C., Clampin, M., et al. 2006, SSRv, 123, 485
Giesers, B., Kamann, S., Dreizler, S., et al. 2019, A\&A, 632, A3
Hajdu, G., Pietrzyński, G., Jurcsik, J., et al. 2021, arXiv:2105.03750
Helmi, A., Babusiaux, C., Koppelman, H. H., et al. 2018, Natur, 563, 85
Helmi, A., White, S. D. M., de Zeeuw, P. T., \& Zhao, H. 1999, Natur, 402, 53
Hobart, M. A., Pena, J. H., \& Peniche, R. 1991, RMxAA, 22, 275
Hubble, E. P. 1926, ApJ, 64, 321
Inno, L., Matsunaga, N., Romaniello, M., et al. 2015, A\&A, 576, A30
Jayasinghe, T., Stanek, K. Z., Kochanek, C. S., et al. 2019, MNRAS, 485, 961
Jones, R. V., Carney, B. W., \& Latham, D. W. 1988a, ApJ, 326, 312
Jones, R. V., Carney, B. W., \& Latham, D. W. 1988b, ApJ, 332, 206
Jones, R. V., Carney, B. W., Latham, D. W., \& Kurucz, R. L. 1987a, ApJ, 314, 605
Jones, R. V., Carney, B. W., Latham, D. W., \& Kurucz, R. L. 1987b, ApJ, 312, 254
Joy, A. H. 1939, ApJ, 89, 356
Keller, S. C., Murphy, S., Prior, S., DaCosta, G., \& Schmidt, B. 2008, ApJ, 678, 851
Kervella, P., Gallenne, A., Evans, N. R., et al. 2019, A\&A, 623, A117
Kinman, T. D., Cacciari, C., Bragaglia, A., Smart, R., \& Spagna, A. 2012, MNRAS, 422, 2116
Kraft, R. P., \& Ivans, I. I. 2003, PASP, 115, 143
Kunder, A., Pérez-Villegas, A., Rich, R. M., et al. 2020, AJ, 159, 270
Layden, A. C. 1994, AJ, 108, 1016
Layden, A. C. 1995, AJ, 110, 2288
Layden, A. C., \& Sarajedini, A. 2003, AJ, 125, 208
Liu, T. 1991, PASP, 103, 205
Liu, T., \& Janes, K. A. 1989, ApJS, 69, 593
Liu, T., \& Janes, K. A. 1990, ApJ, 354, 273
Longmore, A. J., Fernley, J. A., \& Jameson, R. F. 1986, MNRAS, 220, 279
Magurno, D., Sneden, C., Braga, V. F., et al. 2018, ApJ, 864, 57
Majewski, S. R., Schiavon, R. P., Frinchaboy, P. M., et al. 2017, AJ, 154, 94
Massari, D., Koppelman, H. H., \& Helmi, A. 2019, A\&A, 630, L4
Minniti, D., Hempel, M., Toledo, I., et al. 2011, A\&A, 527, A81
Monachesi, A., Gómez, F. A., Grand, R. J. J., et al. 2019, MNRAS, 485, 2589
Moore, C. E. 1972, A Multiplet Table of Astrophysical Interest-Pt.1: Table of Multiplets-Pt.2: Finding List of all Lines in the Table of Multiplets (Washington, DC: US Department of Commerce)
Mucciarelli, A., Lovisi, L., Ferraro, F. R., et al. 2014, ApJ, 797, 43
Mullen, J. P., Marengo, M., Martínez-Vázquez, C. E., et al. 2021, ApJ, 912, 144
Myeong, G. C., Evans, N. W., Belokurov, V., Sanders, J. L., \& Koposov, S. E. 2018, ApJL, 863, L28
Oke, J. B. 1966, ApJ, 145, 468
Oort, J. H. 1927, BAN, 3, 275
Piersimoni, A. M., Bono, G., \& Ripepi, V. 2002, AJ, 124, 1528
Pietrukowicz, P., Kozłowski, S., Skowron, J., et al. 2015, ApJ, 811, 113
Pietrzyński, G., Thompson, I. B., Gieren, W., et al. 2012, Natur, 484, 75
Pojmanski, G. 1997, AcA, 47, 467
Prudil, Z., Hanke, M., Lemasle, B., et al. 2021, arXiv:2102.01090
Riess, A. G., Casertano, S., Yuan, W., Macri, L. M., \& Scolnic, D. 2019, ApJ, 876, 85
Searle, L., \& Zinn, R. 1978, ApJ, 225, 357
Sesar, B. 2012, AJ, 144, 114
Sesar, B., Hernitschek, N., Mitrović, S., et al. 2017, AJ, 153, 204
Sesar, B., Ivezić, Ž., Grammer, S. H., et al. 2010, ApJ, 708, 717
Sesar, B., Jurić, M., \& Ivezić, Z̆. 2011, ApJ, 731, 4
Skillen, I., Fernley, J. A., Stobie, R. S., \& Jameson, R. F. 1993a, MNRAS, 265, 301
Skillen, I., Fernley, J. A., Stobie, R. S., Marang, F., \& Jameson, R. F. 1993b, SAAOC, 15, 90
Smolec, R., Pietrzyński, G., Graczyk, D., et al. 2013, MNRAS, 428, 3034
Sneden, C., Preston, G. W., Chadid, M., \& Adamów, M. 2017, ApJ, 848, 68
Soszyński, I., Udalski, A., Wrona, M., et al. 2019, AcA, 69, 321
Spergel, D., Gehrels, N., Baltay, C., et al. 2015, arXiv:1503.03757

Steinmetz, M., Zwitter, T., Siebert, A., et al. 2006, AJ, 132, 1645
Stetson, P. B., Braga, V. F., Dall’Ora, M., et al. 2014, PASP, 126, 521
Stringer, K. M., Long, J. P., Macri, L. M., et al. 2019, AJ, 158, 16
Su, D. Q., Cui, X., Wang, Y., \& Yao, Z. 1998, Proc. SPIE, 3352, 76
Tamura, N., Takato, N., Shimono, A., et al. 2018, Proc. SPIE, 10702, 107021C
Taylor, W., Cirasuolo, M., Afonso, J., et al. 2018, Proc. SPIE, 10702, 421
Vasiliev, E. 2019, MNRAS, 484, 2832
Vasiliev, E., \& Baumgardt, H. 2021, arXiv:2102.09568

Wan, Z., Oliver, W. H., Baumgardt, H., et al. 2021, MNRAS, 502, 4513
Wesselink, A. J. 1946, BAN, 10, 91
Wright, E. L., Eisenhardt, P. R. M., Mainzer, A. K., et al. 2010, AJ, 140, 1868
Xue, X.-X., Rix, H.-W., Ma, Z., et al. 2015, ApJ, 809, 144
Yanny, B., Rockosi, C., Newberg, H. J., et al. 2009, AJ, 137, 4377
Zinn, R., Chen, X., Layden, A. C., \& Casetti-Dinescu, D. I. 2020, MNRAS, 492, 2161


[^0]:    

    Original content from this work may be used under the terms of the Creative Commons Attribution 4.0 licence. Any further distribution of this work must maintain attribution to the author(s) and the title of the work, journal citation and DOI.

[^1]:    ${ }^{24}$ There is only one exception to this trend: the RRc variable T Sex for which the phase difference is $\sim 0.12$. This large offset might be explained by the fact that this variable is multiperiodic; the shape of its light curve and its luminosity amplitude change night by night (Hobart et al. 1991). It is also worth noting that recent photometric surveys (e.g., ASAS, TESS, Pojmanski 1997; Benkő et al. 2021) display a very narrow light curve for T Sex, meaning that the multiperiodic behavior might have been a transient phenomenon.

